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Analysis of the Wtb vertex from the measurement of triple-differential angular decay rates of single top quarks produced in the t -channel at $\sqrt{s} = 8$ TeV with the ATLAS detector

The ATLAS Collaboration

The electroweak production and subsequent decay of single top quarks in the t -channel is determined by the properties of the Wtb vertex, which can be described by the complex parameters of an effective Lagrangian. An analysis of a triple-differential decay rate in t -channel production is used to simultaneously determine five generalised helicity fractions and phases, as well as the polarisation of the produced top quark. The complex parameters are then constrained. This analysis is based on 20.2 fb^{-1} of proton–proton collision data at a centre-of-mass energy of 8 TeV collected with the ATLAS detector at the LHC. The fraction of decays containing transversely polarised W bosons is measured to be $f_1 = 0.30 \pm 0.05$. The phase between amplitudes for transversely and longitudinally polarised W bosons recoiling against left-handed b -quarks is measured to be $\delta_- = 0.002\pi^{+0.016\pi}_{-0.017\pi}$, giving no indication of CP violation. The fractions of longitudinal or transverse W bosons accompanied by right-handed b -quarks are also constrained. Based on these measurements, limits are placed at 95% CL on the ratio of the complex coupling parameters $\text{Re}[g_R/V_L] \in [-0.12, 0.17]$ and $\text{Im}[g_R/V_L] \in [-0.07, 0.06]$. Constraints are also placed on the ratios $|V_R/V_L|$ and $|g_L/V_L|$. In addition, the polarisation of single top quarks in the t -channel is constrained to be $P > 0.72$ (95% CL). None of the above measurements make assumptions about the value of any of the other parameters or couplings and all of them are in agreement with the Standard Model.

1 Introduction

The top quark is the heaviest known fundamental particle, making the measurement of its production and decay kinematic properties an important probe of physical processes beyond the Standard Model (SM). Within the SM, the top quark decays predominantly through the electroweak interaction to an on-shell W boson and a b -quark. Due to its large mass [1], its lifetime $O(10^{-25} \text{ s})$ is smaller than its hadronisation time-scale $O(10^{-24} \text{ s})$, allowing this quark to be studied as a free quark. Since the top-quark lifetime is also shorter than the depolarisation timescale $O(10^{-21} \text{ s})$ [2] and the W boson is produced on-shell in the top-quark decay, the top-quark spin information is directly transferred to its decay products. Comparing angular measurements of the decay products of polarised top quarks with precise SM predictions provides a unique way to study the non-SM couplings in the Wtb vertex [3]. The normalised triple-differential cross-section (to be defined in Section 2) is the joint probability distribution in all three of the angles determining the kinematics of the decay $t \rightarrow Wb$ from a polarised initial state. Its analysis is the most complete investigation of the dynamics of top-quark decay undertaken to date.

At hadron colliders, top quarks are produced predominantly in pairs ($t\bar{t}$) via the flavour-conserving strong interaction, while an alternative process produces single top quarks through the electroweak interaction. Although the $t\bar{t}$ production cross-section is larger than that of single-top-quark production, top quarks are produced unpolarised because of parity conservation in quantum chromodynamics (QCD) [4], contrary to what happens for single top quarks. At the Large Hadron Collider (LHC) [5], in proton–proton (pp) collision data, the t -channel is the dominant process for producing single top quarks used for the measurements presented in this paper. Figure 1 shows the two representative leading-order (LO) Feynman diagrams for t -channel single-top-quark production. In these two diagrams, a light-flavour quark q (i.e. u - or \bar{d} -quark) from one of the colliding protons interacts with a b -quark by exchanging a virtual W boson, producing a top quark t and a recoiling light-flavour quark q' , called the spectator quark. The b -quark comes either directly from another colliding proton in the five-flavour scheme (5FS) or $2 \rightarrow 2$ process (a) or from a gluon splitting in the four-flavour scheme¹ (4FS) or $2 \rightarrow 3$ process (b). In pp collisions at $\sqrt{s} = 8 \text{ TeV}$, the predicted t -channel production cross-section using the 5FS is $87.8^{+3.4}_{-1.9} \text{ pb}$ [6], calculated at next-to-leading order (NLO) in QCD with resummed next-to-next-to-leading logarithmic (NNLL) accuracy, and called approximate next-to-next-to-leading order (NNLO) in the following. The calculation assumes a top-quark mass of 172.5 GeV and uses the MSTW2008 NNLO [7, 8] parton distribution function (PDF) set. The uncertainties correspond to the sum in quadrature of the uncertainty obtained from the MSTW2008 NNLO PDF set at the 90% confidence level (CL) and the factorisation and renormalisation scale uncertainties.

As a consequence of the vector–axial (V–A) form of the Wtb vertex in the SM, the spin of single top quarks in t -channel production is predominantly aligned along the direction of the spectator-quark momentum [9].

Probes of new physics phenomena affecting the production or decay of the top quark can be parameterised with a series of effective couplings at each vertex [10, 11]; in the t -channel single-top-quark production, both production and decay proceed through the Wtb vertex, and thus are sensitive to the same set of effective couplings.

¹ In the 5FS the b -quarks are treated as massless in the parton distribution functions, while in the 4FS, the parton distribution functions only contain parton distributions for the quarks lighter than the b -quark and b -quarks are treated as massive.

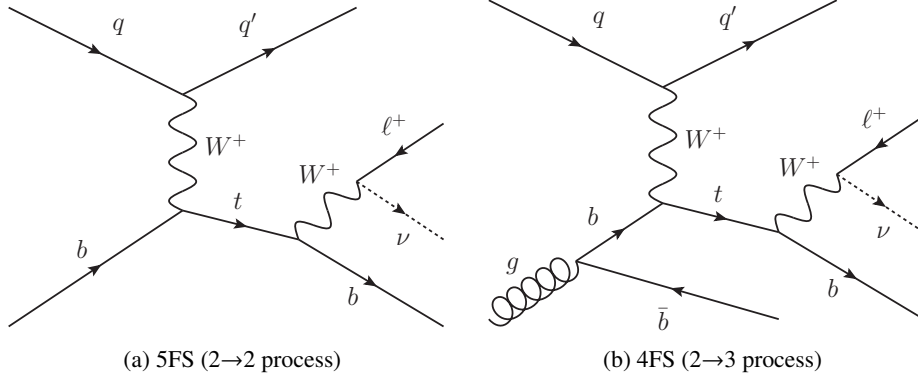


Figure 1: Representative LO Feynman diagrams for t -channel single-top-quark production and decay. Here q represents a u - or \bar{d} -quark, and q' represents (a) a d - or \bar{u} -quark, respectively, in which the initial b -quark arises from a sea b -quark in the 5FS or $2 \rightarrow 2$ process, or (b) a gluon splitting into a $b\bar{b}$ pair in the 4FS or $2 \rightarrow 3$ process.

New physics can be described by an effective Lagrangian, \mathcal{L}_{eff} , represented by dimension-five and dimension-six operators in the framework of effective field theory [12, 13]

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda_{\text{NP}}} \mathcal{L}_5 + \frac{1}{\Lambda_{\text{NP}}^2} \mathcal{L}_6 + \dots,$$

where \mathcal{L}_{SM} represents the SM Lagrangian of dimension four, \mathcal{L}_5 and \mathcal{L}_6 represent the contributions from dimension-five and dimension-six operators invariant under the SM gauge symmetry, and Λ_{NP} is a new physics scale chosen such that higher-dimension operators are sufficiently suppressed by higher powers of Λ_{NP} . Of the standardised set of operators reported in Ref. [12], only four operators, which are dimension six, contribute independently to the Wtb vertex at LO, allowing these terms to be analysed separately from the rest of the full set of possible operators. In a general Lorentz-covariant Lagrangian, expressed by Refs. [10, 11], corrections to the vertex are absorbed into four non-renormalisable effective complex couplings called anomalous couplings:

$$\mathcal{L}_{\text{eff}} = -\frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (V_L P_L + V_R P_R) t W_\mu^- - \frac{g}{\sqrt{2}} \bar{b} \frac{i\sigma^{\mu\nu} q_\nu}{m_W} (g_L P_L + g_R P_R) t W_\mu^- + \text{h.c.},$$

where the four complex effective couplings $V_{L,R}$, $g_{L,R}$ can be identified with the dimension-six operators' Wilson coefficients [14]. Here, g is the weak coupling constant, and m_W and q_ν are the mass and the four-momentum of the W boson. The terms $P_{L,R} \equiv (1 \mp \gamma^5)/2$ are the left- and right-handed projection operators and $\sigma^{\mu\nu} = i[\gamma^\mu, \gamma^\nu]/2$. The terms $V_{L,R}$ and $g_{L,R}$ are the left- and right-handed vector and tensor complex couplings, respectively. In the SM at LO, all coupling constants vanish, except $V_L = V_{tb}$, which is a quark-mixing element in the Cabibbo–Kobayashi–Maskawa (CKM) matrix. Deviations from these values would provide hints of physics beyond the SM, and furthermore, complex values could imply that the top-quark decay has a CP-violating component [15–19].

Indirect constraints on V_L , V_R , g_L , and g_R were obtained [20, 21] from precision measurements of B -meson decays. These results yield constraints in a six-dimensional space of operator coefficients, where four of them correspond to Wtb couplings. Considering one coefficient at a time results in very tight constraints on a particular combination of V_R and g_L , but if several coefficients are allowed to move simultaneously, then individual bounds are not possible. Very tight constraints on CP-violating interactions have been derived from measurements of electric dipole moments [22]. Those constraints also

depend on combinations of couplings, and in a global fit [23], cannot constrain $\text{Im}[g_R]$ better than direct measurements, as are presented here. Measurements of the W boson helicity fractions in top-quark decays [24–28] are sensitive to the magnitude of combinations of anomalous couplings, which are assumed to be purely real, corresponding to the CP-conserving case. These measurements can only place limits on combinations of couplings, and thus the quoted limits on individual couplings depend on the assumptions made about other couplings while V_L is fixed to the SM value of one. More stringent limits are set either in these analyses on $\text{Re}[g_R]$ by considering the measurements of the t -channel single-top-quark production cross-section [29–31] or by performing a global fit considering the most precise measurements of the W boson helicity fractions at the LHC combined with measurements of single-top-quark production cross-sections for different centre-of-mass energies at the LHC and Tevatron [32]. Direct searches for anomalous couplings in t -channel single-top-quark events set limits simultaneously on either both $\text{Re}[g_R/V_L]$ and $\text{Im}[g_R/V_L]$ [33, 34], or on pairs of couplings [35]. In both cases, analyses assume SM values for the other anomalous couplings.

The goal of this analysis is to simultaneously constrain the full space of parameters governing the Wtb vertex using the triple-differential angular decay rate of single top quarks produced in the t -channel as discussed in Section 2, in which the W boson from the top quark subsequently decays leptonically. Conceptually, this is a measurement of each of the anomalous coupling parameters $V_{L,R}$ and $g_{L,R}$ plus the polarisation P of the top quark, with a full covariance matrix; however, any likelihood function derived from the triple-differential decay rate possesses invariances and/or parameter space boundaries lying quite near to the SM point. Therefore, contours are presented instead, with only $\text{Re}[g_R/V_L]$ and $\text{Im}[g_R/V_L]$ showing approximate elliptical contours and therefore admitting point estimation. The anomalous couplings V_R , g_L and g_R are allowed to be complex and the measurements shown require no assumptions to be made regarding the other anomalous couplings. The analysis is carried out in a Fourier-dual space of coefficients in an angular expansion [36, 37]. This method is chosen because it permits an analytic deconvolution of detector effects including both resolution and efficiency, while permitting a simultaneous determination of the real and imaginary parts of all of the anomalous couplings at the Wtb vertex, in addition to the polarisation of the top quark produced in the t -channel.

This paper is organised as follows. Section 2 defines the coordinate system and parameterisation used in the measurement and the triple-differential formalism applied to polarised single top quarks. Section 3 gives a short description of the ATLAS detector, then Section 4 describes the data samples as well as the simulated event samples used to predict properties of the t -channel signal and background processes. Section 5 describes the event reconstruction for the identification of t -channel events, while Section 6 presents the criteria to define the signal region as well as the control and validation regions. The procedures for modelling background processes are reported in Section 7. The event yields and angular distributions comparing the predictions and the observed data are shown in Section 8. Section 9 describes the efficiency, resolution, and background models used to translate the distribution of true t -channel signal events to the distribution of reconstructed signal and background events, and how the parameters of the model are estimated. Section 10 quantifies the sources of uncertainty important in this measurement. Section 11 presents the resulting central value and covariance matrix for the model parameters and the ratios $\text{Re}[g_R/V_L]$ and $\text{Im}[g_R/V_L]$, and the conclusions are given in Section 12.

2 Triple-differential decay rate of polarised single top quarks

An event-specific coordinate system is defined for analysing the decay of the top quark in its rest frame, using the directions of the spectator quark q' that recoils against the top quark, the W boson from the top-quark decay, and the lepton ℓ (e , μ or τ) from the W boson decay, in the final state depicted in Figure 2. The \hat{z} -axis is chosen along the direction of the W boson momentum, \vec{q} , or equivalently along the direction opposite to the b -quark momentum, boosted into the top-quark rest frame, $\hat{z} \equiv \hat{q} = \vec{q}/|\vec{q}|$. The reconstruction of the W boson and top quark is discussed in Section 6. As mentioned before, the spin of single top quarks, \vec{s}_t , in t -channel production is predominantly aligned along the direction of the spectator-quark momentum, \vec{p}_s , in the top-quark rest frame, $\hat{p}_s = \vec{p}_s/|\vec{p}_s|$ [9]. If this quark defines the spin-analysing direction, the degree of polarisation is shown in Refs. [3, 38, 39] to be $P \equiv \hat{p}_s \cdot \vec{s}_t/|\vec{s}_t| \approx 0.9$ at $\sqrt{s} = 8$ TeV for SM couplings. A three-dimensional right-handed coordinate system is defined from the \hat{q} – \hat{p}_s plane and the perpendicular direction, with $\hat{y} = \hat{p}_s \times \hat{q}$ and $\hat{x} = \hat{y} \times \hat{q}$. In this coordinate system, the direction of the lepton momentum, \vec{p}_ℓ^* , in the W boson rest frame, $\hat{p}_\ell^* = \vec{p}_\ell^*/|\vec{p}_\ell^*|$, is specified by the polar angle θ^* and the azimuthal angle ϕ^* . The third angle θ is defined as the angle between \hat{p}_s and \hat{q} . The angle θ^* is the same angle used to measure the W boson helicity fractions in top-quark decays [24–28].

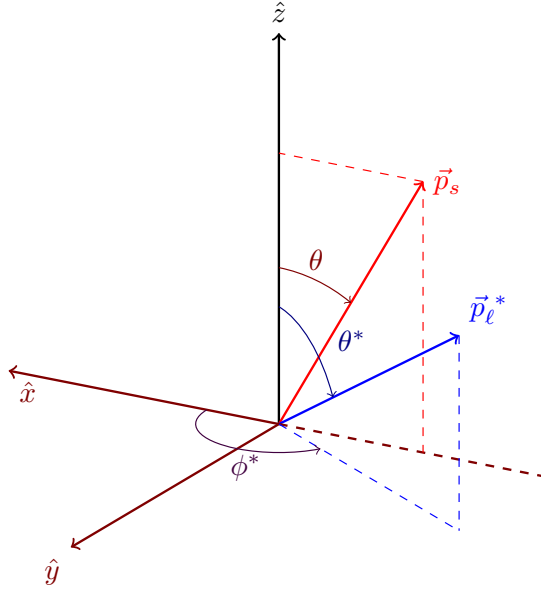


Figure 2: Definition of the right-handed coordinate system with \hat{x} , \hat{y} , and \hat{z} defined as shown from the momentum directions of the W boson, $\hat{q} \equiv \hat{z}$, and the spectator quark, \hat{p}_s with $\hat{y} = \hat{p}_s \times \hat{q}$, in the top-quark rest frame. The angles θ^* and ϕ^* indicate the direction of the lepton momentum, \hat{p}_ℓ^* , while the angle θ indicates the direction of the spectator-quark momentum, \hat{p}_s , in this coordinate system.

These three angles, θ , θ^* , and ϕ^* , arise as a natural choice for measuring a triple-differential distribution for the decay of the top quark, where the W boson subsequently decays leptonically. The $t \rightarrow Wb$ transition is determined by four helicity amplitudes, A_{λ_W, λ_b} , where λ_W and λ_b are the helicities of the W boson and the b -quark, respectively [36]. For $\lambda_b = 1/2$, only the W boson helicities $\lambda_W = 1, 0$ are possible, while for $\lambda_b = -1/2$, $\lambda_W = -1, 0$ are possible. The angular dependence of these transition amplitudes

is given in Ref. [36]. At LO and neglecting the b -quark mass, the helicity amplitudes have a simple dependence on the anomalous couplings. Up to a common proportionality constant, the magnitudes can be expressed as

$$\begin{aligned} |A_{1,\frac{1}{2}}|^2 &\propto 2|x_W V_R - g_L|^2, \\ |A_{0,\frac{1}{2}}|^2 &\propto |V_R - x_W g_L|^2, \\ |A_{-1,-\frac{1}{2}}|^2 &\propto 2|x_W V_L - g_R|^2, \\ |A_{0,-\frac{1}{2}}|^2 &\propto |V_L - x_W g_R|^2, \end{aligned}$$

where $x_W = m_W/m_t$. The relative phases between $A_{1,1/2}$ and $A_{0,1/2}$ and between $A_{-1,-1/2}$ and $A_{0,-1/2}$ are determined by the relative phases between V_R and g_L and between V_L and g_R , respectively.

From the four helicity amplitudes, three fractions can be independently determined. In addition, the interference allows two relative phases between amplitudes to be experimentally determined. These are called the *generalised helicity fractions and phases* [33, 36]:

- f_1 , the fraction of decays containing transversely polarised W bosons,

$$f_1 = \frac{|A_{1,\frac{1}{2}}|^2 + |A_{-1,-\frac{1}{2}}|^2}{|A_{1,\frac{1}{2}}|^2 + |A_{-1,-\frac{1}{2}}|^2 + |A_{0,\frac{1}{2}}|^2 + |A_{0,-\frac{1}{2}}|^2},$$

- f_1^+ , the fraction of b -quarks that are right-handed in events with transversely polarised W bosons,

$$f_1^+ = \frac{|A_{1,\frac{1}{2}}|^2}{|A_{1,\frac{1}{2}}|^2 + |A_{-1,-\frac{1}{2}}|^2},$$

- f_0^+ , the fraction of b -quarks that are right-handed in events with longitudinally polarised W bosons,

$$f_0^+ = \frac{|A_{0,\frac{1}{2}}|^2}{|A_{0,\frac{1}{2}}|^2 + |A_{0,-\frac{1}{2}}|^2},$$

- δ_+ , the phase between amplitudes for longitudinally polarised and transversely polarised W bosons recoiling against right-handed b -quarks,

$$\delta_+ = \arg\left(A_{1,\frac{1}{2}} A_{0,\frac{1}{2}}^*\right),$$

- δ_- , the phase between amplitudes for longitudinally polarised and transversely polarised W bosons recoiling against left-handed b -quarks,

$$\delta_- = \arg\left(A_{-1,-\frac{1}{2}} A_{0,-\frac{1}{2}}^*\right).$$

The fractions f_1 and f_1^+ are related to the quantities F_R , F_0 , and F_L determined by measurements of the W boson helicity fractions in top-quark decays [24–28], with $F_R = f_1 f_1^+$, $F_0 = 1 - f_1$, and $F_L = f_1(1 - f_1^+)$. The fraction f_0^+ is previously unmeasured.

For convenience in what follows, $\vec{\alpha}$ is defined as $\vec{\alpha} \equiv \{f_1, f_1^+, f_0^+, \delta_+, \delta_-\}$. From these five experimental observables, plus the relationships between the helicity amplitudes and the anomalous couplings, one can obtain constraints on all the couplings simultaneously. Additionally, the top-quark polarisation, P , is considered separately from $\vec{\alpha}$ because it depends on the production of the top quark, rather than on its decay.

At LO, the helicity amplitudes, and hence $\vec{\alpha}$ can be expressed as functions of the couplings and the parton masses [19, 40]. Using SM couplings and $m_b = 4.95$ GeV, $m_t = 172.5$ GeV, and $m_W = 80.399$ GeV with the derived analytic expressions for $\vec{\alpha}$, the expected values are

$$f_1 = 0.304, \quad f_1^+ = 0.001, \quad f_0^+ = 6 \cdot 10^{-5}, \quad \delta_+ = \delta_- = 0.0.$$

Calculations at NNLO [41] predict $f_1 = 0.311 \pm 0.005$, and $f_1^+ = 0.0054 \pm 0.0003$, where the largest part of the uncertainty in f_1 comes from the experimental uncertainty of the top-quark mass, while for f_1^+ it arises from uncertainties in α_s and the b -quark mass. An NNLO prediction does not yet exist for f_0^+ , but NLO calculations [40] yield a value < 0.001 .

In Refs. [36, 37] it is shown that the Jacob–Wick helicity formalism [42, 43] applied to the decay of polarised top quarks in t -channel production leads to the following expression for the triple-differential decay rate for polarised top quarks in terms of the three angles (θ , θ^* , and ϕ^*) and the top-quark polarisation,

$$\begin{aligned} \varrho(\theta, \theta^*, \phi^*; P) &= \frac{1}{N} \frac{d^3 N}{d(\cos \theta) d\Omega^*} = \frac{1}{8\pi} \left\{ \frac{3}{4} \left| A_{1, \frac{1}{2}} \right|^2 (1 + P \cos \theta)(1 + \cos \theta^*)^2 \right. \\ &+ \frac{3}{4} \left| A_{-1, -\frac{1}{2}} \right|^2 (1 - P \cos \theta)(1 - \cos \theta^*)^2 \\ &+ \frac{3}{2} \left(\left| A_{0, \frac{1}{2}} \right|^2 (1 - P \cos \theta) + \left| A_{0, -\frac{1}{2}} \right|^2 (1 + P \cos \theta) \right) \sin^2 \theta^* \\ &- \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 + \cos \theta^*) \operatorname{Re} \left[e^{i\phi^*} A_{1, \frac{1}{2}} A_{0, \frac{1}{2}}^* \right] \\ &- \left. \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 - \cos \theta^*) \operatorname{Re} \left[e^{-i\phi^*} A_{-1, -\frac{1}{2}} A_{0, -\frac{1}{2}}^* \right] \right\} \\ &= \sum_{k=0}^1 \sum_{l=0}^2 \sum_{m=-k}^k a_{k,l,m} M_{k,l}^m(\theta, \theta^*, \phi^*), \end{aligned} \quad (1)$$

where $d\Omega^* \equiv d(\cos \theta^*) d\phi^*$ (see Figure 2). The $a_{k,l,m}$ represent the angular coefficients to be determined and $M_{k,l}^m(\theta, \theta^*, \phi^*)$ are orthonormal functions over the three angles defined by the product of two spherical harmonics, $Y_k^m(\theta, 0)$ and $Y_l^m(\theta^*, \phi^*)$,

$$M_{k,l}^m(\theta, \theta^*, \phi^*) = \sqrt{2\pi} Y_k^m(\theta, 0) Y_l^m(\theta^*, \phi^*).$$

The properties of these M -functions are detailed in Ref. [37]. The restriction to $k \leq 1$ and $l \leq 2$ in Eq. (1) is caused by the allowed spin states of the initial- and final-state fermions and the vector boson at the weak vertex.

Only nine of the angular coefficients $a_{k,l,m}$, not taking into account $a_{0,0,0}$, which is constrained by normalisation ($|A_{1,1/2}|^2 + |A_{0,1/2}|^2 + |A_{-1,-1/2}|^2 + |A_{0,-1/2}|^2 = 1$), are non-zero and can be parameterised in terms of the generalised helicity fractions and phases.

The non-zero angular coefficients $a_{k,l,m}(\vec{\alpha}; P)$ are:

$$\begin{aligned}
a_{0,0,0} &= \frac{1}{\sqrt{8\pi}}, \\
a_{0,1,0} &= \frac{\sqrt{3}}{\sqrt{8\pi}} f_1 \left(f_1^+ - \frac{1}{2} \right), \\
a_{0,2,0} &= \frac{1}{\sqrt{40\pi}} \left(\frac{3}{2} f_1 - 1 \right), \\
a_{1,0,0} &= +P \frac{1}{\sqrt{24\pi}} \left(f_1(2f_1^+ - 1) + (1 - f_1)(1 - 2f_0^+) \right), \\
a_{1,1,0} &= +P \frac{1}{\sqrt{32\pi}} f_1, \\
a_{1,2,0} &= +P \frac{1}{\sqrt{480\pi}} \left(f_1(2f_1^+ - 1) - 2(1 - f_1)(1 - 2f_0^+) \right), \\
a_{1,1,1} &= (a_{1,1,-1})^* = -P \frac{1}{\sqrt{16\pi}} \sqrt{f_1(1 - f_1)} \left\{ \sqrt{f_1^+ f_0^+} e^{i\delta_+} + \sqrt{(1 - f_1^+)(1 - f_0^+)} e^{-i\delta_-} \right\}, \\
a_{1,2,1} &= (a_{1,2,-1})^* = -P \frac{1}{\sqrt{80\pi}} \sqrt{f_1(1 - f_1)} \left\{ \sqrt{f_1^+ f_0^+} e^{i\delta_+} - \sqrt{(1 - f_1^+)(1 - f_0^+)} e^{-i\delta_-} \right\}, \quad (2)
\end{aligned}$$

where $(a_{k,l,m})^*$ represents a complex conjugate. All the other angular coefficients are zero in top-quark decays.

Coefficients of M -functions can also be determined from data. In Section 9, techniques are discussed for measuring those coefficients, how to deconvolve them to obtain the coefficients presented here, and hence the parameters $\vec{\alpha}$ and P .

3 ATLAS detector

The ATLAS detector [44] consists of a set of sub-detector systems, cylindrical in the central region and planar in the two endcap regions, that covers almost the full solid angle around the interaction point (IP).² ATLAS is composed of an inner detector (ID) for tracking close to the IP, surrounded by a superconducting solenoid providing a 2 T axial magnetic field, electromagnetic (EM) and hadronic calorimeters, and a muon spectrometer (MS). The ID consists of a silicon pixel detector, a silicon micro-strip detector, providing tracking information within pseudorapidity $|\eta| < 2.5$, and a straw-tube transition radiation tracker that covers $|\eta| < 2.0$. The central EM calorimeter is a lead and liquid-argon (LAr) sampling calorimeter with high granularity, and is divided into a barrel region that covers $|\eta| < 1.5$ and endcap regions

² ATLAS uses a right-handed coordinate system with its origin at the nominal IP in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. The transverse momentum and energy are defined as $p_T = p \sin \theta$ and $E_T = E \sin \theta$, respectively. The ΔR is the distance defined as $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

that cover $1.4 < |\eta| < 3.2$. A steel/scintillator tile calorimeter provides hadronic energy measurements in the central range of $|\eta| < 1.7$. The endcap ($1.5 < |\eta| < 3.2$) and forward regions ($3.1 < |\eta| < 4.9$) are instrumented with LAr calorimeters for both the EM and hadronic energy measurements. The MS consists of three large superconducting toroid magnets with eight coils each, a system of trigger chambers covering $|\eta| < 2.4$, and precision tracking chambers covering $|\eta| < 2.7$. The ATLAS detector employs a three-level trigger system [45], used to select events to be recorded for offline analysis. The first-level trigger is hardware-based, implemented in custom-built electronics and it uses a subset of the detector information to reduce the physical event rate from 40 MHz to at most 75 kHz. The second-level trigger and the final event filter, collectively referred to as the high-level trigger (HLT), are software-based and together reduce the event rate to about 400 Hz.

4 Data and simulation samples

The analysis is performed using data from pp collisions delivered by the LHC in 2012 at $\sqrt{s} = 8$ TeV and recorded by the ATLAS detector. Stringent detector and data quality requirements were applied, resulting in a data sample corresponding to a total integrated luminosity of 20.2 fb^{-1} [46]. The events were selected by single-lepton³ triggers [45, 47], imposing at the HLT a threshold of 24 GeV on the transverse energy (E_T) of electrons and on the transverse momentum (p_T) of muons, along with isolation requirements. To recover efficiency for higher- p_T leptons, the isolated lepton triggers were complemented by triggers without isolation requirements, but with a threshold raised to 60 GeV for electrons and to 36 GeV for muons.

Samples of events generated using Monte Carlo (MC) simulations were produced using different event generators interfaced to various parton showering (PS) and hadronisation generators. Minimum-bias events simulated with the PYTHIA8 generator (ver. 8.1) [48] were overlaid to model the effect of multiple pp collisions per bunch crossing (pile-up). The distribution of the average number of pile-up interactions in the simulation is reweighted to match the corresponding distribution in data, which has an average of 21 [46]. The events were processed using the same reconstruction and analysis chain as for data events.

Single-top-quark t -channel events were generated with the NLO POWHEG-Box generator (rev. 2556) [49] with the CT10f4 [50] PDF set, using the 4FS for the matrix-element (ME) calculations [51]. The renormalisation and factorisation scales were set to $\mu_R^2 = \mu_F^2 = 16(m_b^2 + p_{T,b}^2)$, where m_b is the mass of the b -quark and $p_{T,b}$ is the transverse momentum of the b -quark from the initial gluon splitting. Top quarks were decayed using MADSPIN [52], which preserves all spin correlations. Additional t -channel samples were produced with the LO PROTOS generator (ver. 2.2b) [53] using the CTEQ6L1 PDF set [54] within the 4FS. Thus in addition to a SM sample, samples with anomalous couplings enabled in both the production and the decay vertices were produced using the PROTOS generator, varying simultaneously V_L with either $\text{Re}[V_R] \in [0.25, 0.50]$, $\text{Re}[g_R] \in [-0.26, 0.18]$ or $\text{Im}[g_R] \in [-0.23, 0.23]$, such that the top-quark width was invariant. The factorisation scale was set to $\mu_F^2 = -p_W^2$ for the spectator quark and $\mu_F^2 = p_{\bar{b}}^2 + m_{\bar{b}}^2$ for the gluon, where p_W and $p_{\bar{b}}$ are the three-momenta of the exchanged W boson and of the \bar{b} -quark originating from the gluon splitting (the spectator \bar{b} -quark), respectively. In order to compare different LO generators, another sample of signal events was produced with the multi-leg LO ACERMC generator (ver. 3.8) [55] using the CTEQ6L1 PDF set. This generator incorporates both 4FS and 5FS, featuring an automated procedure to remove the overlap in phase space between the two schemes [56]. The factorisation and renormalisation scales were set to $\mu_F = \mu_R = m_t = 172.5$ GeV.

³ Henceforth, “lepton” indicates electron or muon, and does not include τ leptons.

In this analysis, all simulated signal event samples are normalised using the production cross-section mentioned in Section 1. Simulation samples produced with POWHEG-Box are used for predicting the acceptance and the template shape of the t -channel signal. To estimate the efficiency and resolution models, the simulation samples in which parton-level information is well defined, i.e. those produced with either PROTONS or ACERMC, are used.

Samples of simulated events for $t\bar{t}$ production and electroweak production of single top quarks in the associated Wt and s -channel were produced using the NLO POWHEG-Box generator (rev. 2819, rev. 3026) coupled with the CT10 [50] PDF set. The t - and s -channel processes do not interfere even at NLO in QCD and are thus well defined with that precision [57]. For Wt associated production, the diagram removal scheme is used to eliminate overlaps between this process and $t\bar{t}$ production at NLO. In the $t\bar{t}$ sample, the resummation damping factor⁴ h_{damp} was set to the top-quark mass [58]. An additional $t\bar{t}$ sample with anomalous couplings enabled in the decay vertex was produced using the PROTONS generator (ver. 2.2) coupled with the CTEQ6L1 PDF set. This sample is used to take into account the dependence of $t\bar{t}$ background upon the value of the anomalous couplings.

For all simulated event samples mentioned above, the PS, hadronisation and underlying event (UE) were added using PYTHIA (ver. 6.426, ver. 6.427) [59] with the Perugia 2011C set of tuned parameters (P2011C tune) [60] and the CTEQ6L1 PDF set. The TAUOLA [61] program and the PHOTOS [62] algorithm were used to properly simulate decays of polarised τ leptons including spin correlations and to generate quantum electrodynamics (QED) radiative corrections in decays to account for photon radiation. All these processes were simulated assuming a top-quark mass of 172.5 GeV, and the decay of the top quark was assumed to be 100% $t \rightarrow Wb$.

For estimating the t -channel and $t\bar{t}$ generator modelling uncertainties, additional samples were produced using alternative generators or parameter variations. For studying the top-quark mass dependence, supplementary single-top-quark and $t\bar{t}$ simulated event samples with different top-quark masses were generated. These topics are further discussed in Section 10 and Section 11, respectively.

Vector-boson production in association with jets was simulated using the multi-leg LO SHERPA generator (ver. 1.4.1) [63] with its own parameter tune and the CT10 PDF set. Thus, W +jets and Z +jets events with up to four additional partons were generated and the contributions of W/Z +light-jets and W/Z +heavy-jets ($W/Z+bb$, $W/Z+cc$, $W/Z+c$) were simulated separately. SHERPA was also used to generate the hard process, but also for the PS, hadronisation and the UE, using the CKKW method [64] to remove overlaps between the partonic configurations generated by the ME and by the PS. Samples of diboson events (WW , WZ , and ZZ), containing up to three additional partons where at least one of the bosons decays leptonically, were also produced using the SHERPA generator (ver. 1.4.1) with the CT10 PDF set.

All baseline simulated event samples were passed through the full simulation of the ATLAS detector [65] based on the GEANT4 framework [66] while PROTONS simulated event samples and alternative samples used to estimate systematic uncertainties were processed through a faster simulation using the ATLFAT2 framework [67].

⁴ The resummation damping factor, h_{damp} , is one of the parameters controlling the ME/PS matching in POWHEG and effectively regulates the high- p_T gluon radiation. In the used POWHEG-Box revision, $h_{\text{damp}} = \infty$ was the default value.

5 Event reconstruction

Electron candidates are reconstructed from isolated energy deposits in the EM calorimeter associated with ID tracks fulfilling strict quality requirements [68]. These electrons are required to satisfy $E_T = E_{\text{cluster}} / \sin(\theta_{\text{track}}) > 25$ GeV and $|\eta_{\text{cluster}}| < 2.47$, where E_{cluster} and η_{cluster} denote the energy and the pseudorapidity of the cluster of energy deposits in the EM calorimeter, and θ_{track} denotes the polar angle of the ID track associated with this cluster. Clusters in the EM calorimeter barrel–endcap transition region, corresponding to $1.37 < |\eta_{\text{cluster}}| < 1.52$, are excluded. Muon candidates are reconstructed using combined information from the ID tracks and the MS [69]. They are required to have $p_T > 25$ GeV and $|\eta| < 2.5$. The electron and muon candidates must fulfil additional isolation requirements, as described in Ref. [70], in order to reduce contributions from misidentified jets, non-prompt leptons from the decay of heavy-flavour quarks and non-prompt electrons from photon conversions.

Jets are reconstructed using the anti- k_t algorithm [71, 72] with a radius parameter of 0.4, using topological clusters of calorimeter energy deposits [73] as inputs to the jet finding. The clusters are calibrated with a local cluster weighting method [73]. The jet energy is further corrected for the effect of multiple pp interactions. Jets are calibrated using an energy- and η -dependent simulation-based scheme, with in situ corrections based on data [74]. To reject jets from pile-up events, a so-called jet-vertex-fraction (JVF) criterion [75] is applied to the jets with $p_T < 50$ GeV and $|\eta| < 2.4$: at least 50% of the scalar sum of the p_T of the tracks associated with a jet is required to be from tracks compatible with the primary vertex.⁵ Only events containing reconstructed jets with $p_T > 30$ GeV and $|\eta| < 4.5$ are considered. The p_T threshold is raised to 35 GeV for the jets in the calorimeter endcap–forward transition region, corresponding to $2.7 < |\eta| < 3.5$ [29]. Jets identified as likely to contain b -hadrons are tagged as b -jets. The b -tagging is performed using a neural network (NN) which combines three different algorithms exploiting the properties of a b -hadron decay in a jet [76]. The b -tagging algorithm, only applied to jets within the coverage of the ID (i.e. $|\eta| < 2.5$), is optimised to improve the rejection of c -quark jets, since W boson production in association with c -quarks is a major background for the selected final state. The requirement applied to the NN discriminant corresponds to a b -tagging efficiency of 50%, with mis-tagging rates of 3.9% and 0.07% for c -quark jets and light-flavour jets (u -, d -, s -quark or gluon g), respectively, as predicted in simulated $t\bar{t}$ events and calibrated with data [77, 78].

The missing transverse momentum, with magnitude E_T^{miss} , is reconstructed from the vector sum of energy deposits in the calorimeter projected onto the transverse plane [79]. The energies of all clusters are corrected using the local cluster weighting method. Clusters associated with high- p_T jets and electrons are further calibrated using their respective energy corrections. In addition, contributions from the p_T of the selected muons are also included in the calculation. The E_T^{miss} is taken as a measurement of the undetectable particles, and is affected by energy losses due to detector inefficiencies and acceptance, and by energy resolution.

⁵ A primary-vertex candidate is defined as a reconstructed vertex with at least five associated tracks with $p_T > 400$ MeV. The primary vertex associated with the hard-scattering collision is the candidate with the largest sum of the squared p_T of the associated tracks.

6 Event selection in the signal, control, and validation regions

The signal event candidates are selected by requiring a single prompt isolated lepton,⁶ significant E_T^{miss} , and exactly two jets. All these objects must satisfy the criteria described in Section 5, and the E_T^{miss} is required to be larger than 30 GeV. One of the jets must be identified as a b -tagged jet with $|\eta| < 2.5$ while the second jet, also called the spectator jet, is required to be untagged and produced in the forward direction. Events containing additional jets are vetoed to suppress background from $t\bar{t}$ production. The spectator \bar{b} -quark originating from the gluon splitting (4FS), as shown in Figure 1(b), can result in an additional b -tagged jet. This jet is expected to have a softer p_T spectrum and a broader η distribution than the b -tagged jet produced in the top-quark decay. It is generally not detected in the experiment and these events pass the event selection. Events are required to contain at least one good primary vertex candidate, and no jets failing to satisfy reconstruction quality criteria. In addition, the transverse mass of the lepton– E_T^{miss} system,

$$m_T(\ell E_T^{\text{miss}}) = \sqrt{2p_T(\ell) \cdot E_T^{\text{miss}} \left[1 - \cos(\Delta\phi(\ell, E_T^{\text{miss}})) \right]},$$

where $\Delta\phi(\ell, E_T^{\text{miss}})$ is the difference in azimuthal angle between the lepton momentum and the E_T^{miss} direction, is required to be larger than 50 GeV in order to reduce the multijet background contribution. Further reduction of this background is achieved by imposing a requirement on the lepton p_T to events in which the lepton and leading jet (j_1) are back-to-back [29, 33, 80],

$$p_T(\ell) > 40 \left(\frac{|\Delta\phi(j_1, \ell)| - 1}{\pi - 1} \right) \text{ GeV},$$

where $\Delta\phi(j_1, \ell)$ is the difference in azimuthal angle between the lepton momentum and the leading jet. To reduce the dilepton backgrounds, events containing an additional lepton, identified with less stringent criteria (referred to as a loose lepton) and with a p_T threshold lowered to 10 GeV, are rejected. Finally, two additional requirements are applied in order to remove a mis-modelling between data and prediction seen in the W +jets control and validation regions, in the $|\eta|$ distribution of the non- b -jet and in the $|\Delta\eta|$ distribution between the two required jets: $|\eta(\text{non-}b\text{-jet})| < 3.6$ and $|\Delta\eta(\text{non-}b\text{-jet}, b\text{-jet})| < 4.5$.

The W boson originating from the decay of the top quark is reconstructed from the momenta of the lepton and the neutrino by imposing four-momentum conservation. Since the neutrino escapes undetected, the x and y components of the reconstructed E_T^{miss} are assumed to correspond to the p_T of the neutrino. The unmeasured longitudinal component of the neutrino momentum, p_v^z , is computed by imposing a W boson mass constraint on the lepton–neutrino system. A quadratic expression is found for p_v^z . If there are two real solutions, the solution closer to zero is taken. If the solutions are complex, the assumption of the neutrino being the only contributor to the E_T^{miss} is not valid.⁷ Therefore, the reconstructed E_T^{miss} is rescaled, preserving its direction, in order to have physical (real) solutions for p_v^z . This generally results in two solutions for the rescaled E_T^{miss} . If just one solution of the rescaled E_T^{miss} is positive, this is chosen. If both are positive, the one closer to the initial E_T^{miss} is chosen. The top-quark candidate is then reconstructed by combining the four-momenta of the reconstructed W boson and the selected b -tagged jet. Finally, the momenta of the W boson and spectator jet are boosted into the top-quark rest frame to

⁶ This analysis considers only W boson decay modes to an electron or a muon. Events in which the W boson decays to a τ lepton are included if the τ subsequently decays to an electron or a muon.

⁷ Although it is true that at LO the neutrino is the main contributor to the E_T^{miss} , there may be other contributors, such as extra neutrinos (from b -hadron and τ decays), additional p_T contributions (initial/final-state radiation effects), miscalibration of E_T^{miss} , fake E_T^{miss} due to the detector energy resolution and acceptance.

obtain \vec{q} and \vec{p}_s , used to define the coordinate system in Figure 2, and the lepton is boosted into the W boson rest frame to obtain \vec{p}_ℓ^* .

In addition to this basic event selection, which defines the preselected region, further discrimination between the t -channel signal events and background events is achieved by applying additional criteria:

- The pseudorapidity of the non- b -tagged jet must satisfy $|\eta(\text{non-}b\text{-jet})| > 2.0$, since the spectator jet tends to be produced in the forward region in the t -channel signature.
- The scalar sum of the p_T of all final-state objects (lepton, jets and E_T^{miss}), H_T , must be larger than 195 GeV, since the H_T distributions of the backgrounds peak at lower values (in particular for the W +jets contribution) than the t -channel signature.
- The mass of the top quark reconstructed from its decay products, $m(\ell\nu b)$, is required to be within 130–200 GeV, to reject background events from processes not involving top quarks.
- The absolute difference in η between the non- b -tagged jet and the b -jet, $|\Delta\eta(\text{non-}b\text{-jet}, b\text{-jet})|$, must be larger than 1.5, to further reduce $t\bar{t}$ contributions.

These criteria are based on the selection requirements used in Ref. [33], re-optimised using MC simulation at $\sqrt{s} = 8$ TeV [34]. Thus, these criteria together with the signal preselection define the signal region of this analysis.

The distributions of the four variables used to define the signal region are shown in Figure 3 at the preselection stage. The simulated signal and background distributions are scaled to their theoretical predictions except the multijet background, which is estimated using data-driven techniques described in Section 7. The W +jets, top-quark backgrounds and t -channel distributions are normalised to the results of the maximum-likelihood fit, also described in Section 7. In Figure 3(a), the well-modelled bump around $|\eta| = 2.5$ is due to a combination of the JVF requirement, which is applied to jets with $p_T < 50$ GeV and $|\eta| < 2.4$, and the increased p_T requirement on jets in the calorimeter endcap–forward transition region ($2.7 < |\eta| < 3.5$). These two requirements are described in Section 5.

To estimate the rates and validate the modelling of the dominant background contributions, the simulated events are compared to the data in three dedicated background-enriched regions:

- A control region dominated by $t\bar{t}$ events is defined by considering preselected events containing two additional non- b -tagged jets (i.e. four jets are required since just one of them is required to be b -tagged).
- A control region enriched in W +jets events, and dominated by W +heavy-jets, is defined in order to control the modelling of the background. The events selected in this control region are the ones satisfying the preselection criteria and failing to satisfy any of the four requirements in the selection criteria. The flavour composition of this control region is similar to that of the signal region.
- A third region is defined as a validation region dominated by W +jets events to further control the modelling of the shapes of the W +jets background. Events in this validation region are selected by considering the preselection criteria with a relaxed b -tagging efficiency requirement of 80%. In addition, all events satisfying the tighter signal b -tagging efficiency requirement of 50% are excluded. This region has much larger enrichment in W +jets events although the flavour composition differs from that of the signal region.

The two control regions are used to extract the normalisation of $t\bar{t}$ and W +jets as described in Section 7.

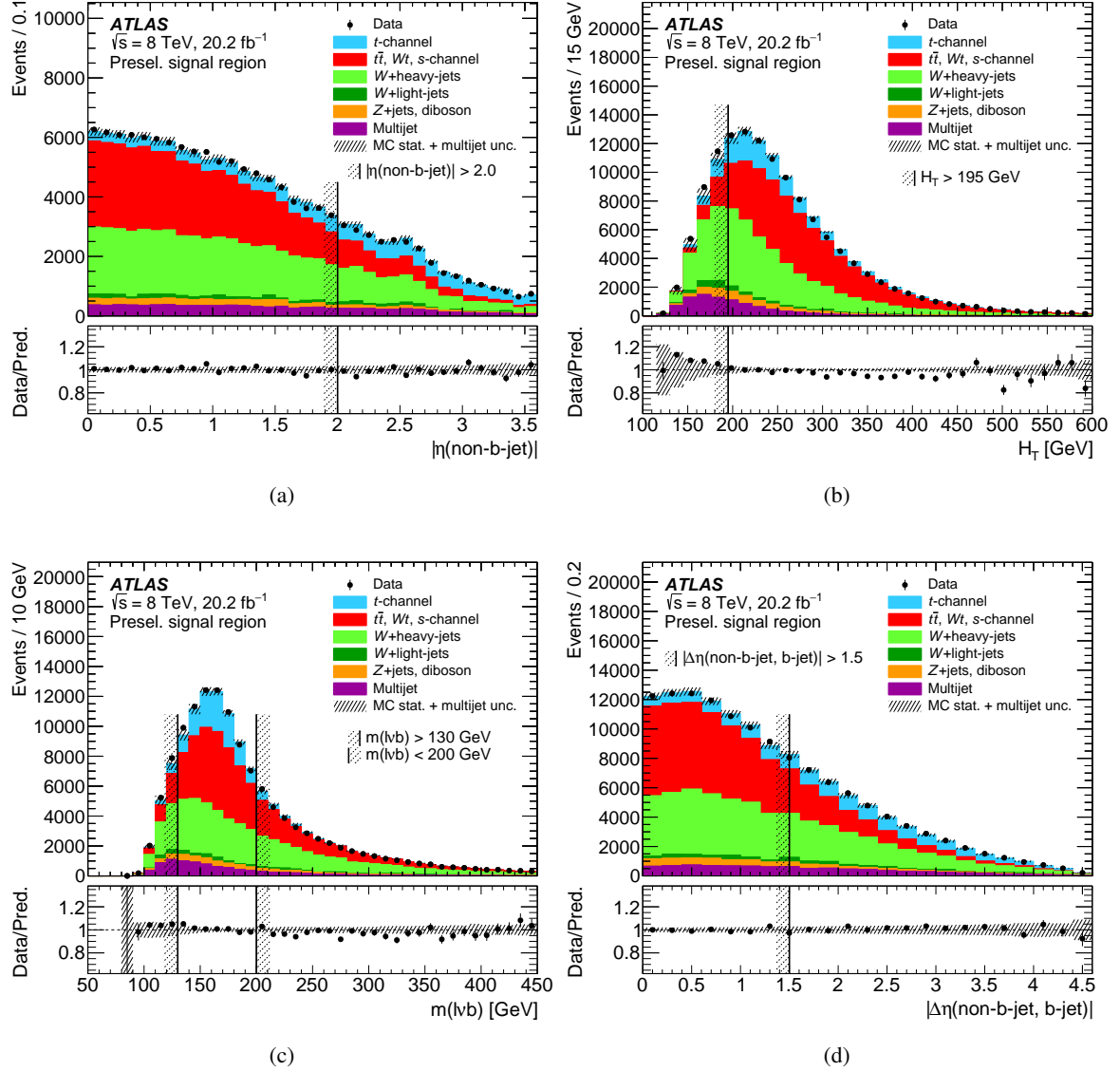


Figure 3: Distributions of (a) $|\eta(\text{non-}b\text{-jet})|$, (b) the scalar sum of the p_T of all final-state objects, H_T , (c) reconstructed top-quark mass, $m(\ell vb)$, and (d) $|\Delta\eta(\text{non-}b\text{-jet}, b\text{-jet})|$ in the signal preselected region for the electron and muon channels merged. The prediction is compared to data, shown as the black points with statistical uncertainties. The multijet background is estimated using data-driven techniques, while contributions from simulated W +jets, top-quark backgrounds and t -channel event samples are normalised to the results of a maximum-likelihood fit to event yields in the signal and control regions. The uncertainty bands correspond to the uncertainties due to the size of the simulated event samples added in quadrature with the data-driven normalisation uncertainty of 70% estimated for the multijet contribution. The lower plots show the ratio of data to prediction in each bin. The regions excluded by the selection criteria are shown by vertical black lines and dashed areas.

7 Background estimation and normalisation

The largest background contributions to single-top-quark t -channel production arise from $t\bar{t}$ and W +jets production. The former is difficult to distinguish from the signal since $t\bar{t}$ events contain real top quarks in the final state. The W +jets production contributes to the background if there is a b -quark in the final state or due to mis-tagging of jets containing other quark flavours. Multijet production via the strong interaction can contribute as well if, in addition to two reconstructed jets, an extra jet is misidentified as an isolated lepton, or if a non-prompt lepton appears to be isolated (both referred to as fake leptons). Other minor backgrounds originate from single-top-quark Wt -channel and s -channel, Z +jets and diboson production.

For all background processes, except multijet production, the normalisation is initially estimated by using the MC simulation scaled with the theoretical cross-section prediction, and the event distribution modelling is taken from simulation.

The $t\bar{t}$ events are normalised with the $t\bar{t}$ production cross-section calculated at NNLO in QCD including resummation of NNLL soft gluon terms with Top++2.0 [81–86]. Its predicted value is 253^{+13}_{-15} pb calculated according to Ref. [86]. The quoted uncertainty, evaluated according to the PDF4LHC prescription [87], corresponds to the sum in quadrature of the α_s uncertainty and the PDF uncertainty, calculated from the envelope of the uncertainties at 68% CL of the MSTW2008 NNLO, CT10 NNLO [88] and NNPDF2.3 5f FFN [89] PDF sets. The associated Wt -channel events are normalised with the predicted NNLO production cross-section of 22.4 ± 1.5 pb [90] and the s -channel production to the predicted NNLO cross-section of 5.61 ± 0.22 pb [91]. The uncertainties correspond to the sum in quadrature of the uncertainty derived from the MSTW2008 NNLO PDF set at 90% CL and the scale uncertainties.

The inclusive cross-sections of vector-boson production are calculated to NNLO with the FEWZ program [92] and the MSTW2008 NNLO PDF set, with a theoretical uncertainty of 4% and 5% for W +jets and Z +jets, respectively. The cross-sections of diboson processes are calculated at NLO using the MCFM program [93], with a theoretical uncertainty of 5%. For these three background processes the normalisation uncertainty is 34% each. This is the result of adding in quadrature their theory uncertainty and 24% per additional jet, accordingly to the Berends–Giele scaling [94].

The normalisation as well as the event modelling of the multijet background is estimated from data using a matrix method [70, 95]. This method allows the derivation of the true composition of the data sample in terms of prompt (real) and fake leptons from its observed composition in terms of tight (signal selection) and loose leptons. An alternative normalisation and modelling based on the mixed data–simulation jet–electron method [29, 70, 96] and the purely data-driven anti-muon selection [70] are also considered. From the comparison of these two models with the results obtained using the matrix method, an overall normalisation uncertainty of 70% is assigned to the multijet contribution, irrespective of lepton flavour, as done in Ref. [34].

The final t -channel, W +jets and top-quark background ($t\bar{t}$, associated Wt and s -channel) normalisations are estimated through a simultaneous maximum-likelihood fit to the numbers of data events observed in the signal region and the $t\bar{t}$ and W +jets control regions, described in Section 6. The likelihood function [96] is given by the product of Poisson probability terms associated with the fitted regions, combined with the product of Gaussian priors to constrain the background rates to their predictions within the associated uncertainties. In the fit, the t -channel contribution, estimated using POWHEG-Box, is treated as unconstrained. The top-quark background contributions are merged with their relative fractions taken from simulation, and the applied constraint, 6%, is derived from the combination in quadrature of their

cross-section uncertainties. The W +jets contribution is constrained to the normalisation uncertainty of 34% and its flavour composition is taken from simulation. In these three fitted regions the production of a W boson in association with heavy-flavour jets is the dominant contribution to the W +jets background, predicted to be around 95% in each region. The Z +jets and diboson contributions, which are very low in the signal region (2% of the expected total), are merged and fixed to the predictions. The multijet contribution is kept fixed to its data-driven estimate. The overall normalisation scale factors obtained from the maximum-likelihood fit together with the statistical post-fit uncertainties are found to be 1.010 ± 0.005 and 1.128 ± 0.013 for the top-quark and W +jets background contributions, respectively, and 0.909 ± 0.022 for the t -channel signal. The impact on the analysis due to the deviation of these scale factors from unity is negligible and it is taken into account through the W +jets normalisation uncertainty as discussed in Section 10. In the case of the W +jets validation region, used to validate the shapes of the predicted templates, just an overall scale factor for the W +jets component is estimated. It is extracted by matching the total predicted event yields to the number of events observed in this validation region. The results are found to be stable when the prior constraints on the top-quark and W +jets backgrounds are relaxed to 100% of their predicted cross-section in the signal and control regions.

The overall normalisation scale factors are used to control the modelling of the kinematic and angular variable distributions in the signal, control, and validation regions. In the subsequent steps of the analysis, the overall scaling of the t -channel prediction is not relevant, since it is taken from background-subtracted data, while the W +jets and top-quark backgrounds are normalised using these overall scale factors.

8 Event yields and kinematic distributions

Table 1 provides the predicted signal and background event yields for the electron and muon channels merged together in the signal, control, and validation regions after scaling to the results of the maximum-likelihood fit to the data. Observed data yields are also shown. The signal-to-background (S/B) ratio is 0.97 in the signal region while $\lesssim 0.1$ in the control and validation regions.

Process	Signal region	$t\bar{t}$ control region	W +jets control region	W +jets validation region
t -channel	4395 ± 17	1688 ± 12	11601 ± 29	9306 ± 27
$t\bar{t}$, Wt , s -channel	2017 ± 15	62864 ± 77	48120 ± 82	23937 ± 61
W +heavy-jets	1910 ± 49	6898 ± 65	45410 ± 200	157260 ± 480
W +light-jets	87 ± 31	218 ± 38	3110 ± 200	130900 ± 1000
Z +jets, diboson	157 ± 7	1118 ± 37	4734 ± 77	17750 ± 300
Multijet	375 ± 13	862 ± 27	8910 ± 61	20140 ± 120
Total expected	8941 ± 64	73650 ± 120	121890 ± 310	359300 ± 1200
Data	8939	73662	121913	359320
S/B	0.97	0.02	0.11	0.03

Table 1: Predicted and observed data event yields are shown for the merged electron and muon channels in the signal, $t\bar{t}$ and W +jets control and validation regions. The multijet background is estimated using data-driven techniques, while contributions from simulated W +jets, top-quark backgrounds and t -channel event samples are normalised to the results of a maximum-likelihood fit to event yields in the signal and control regions. The uncertainties shown are statistical only. Individual predictions are rounded to two significant digits of the uncertainty while “Total expected” corresponds to the rounding of the sum of full-precision individual predictions. The expected S/B ratios are also given.

Figures 4 and 5 show the distributions of the relevant kinematic distributions used to define the signal region in the $t\bar{t}$ and W +jets control regions while Figure 6 shows the same distributions in the W +jets validation region. Good overall data-to-prediction agreement is found within the uncertainty band shown in these distributions, which only includes the uncertainty due to the size of the simulation samples and the uncertainty in the normalisation of the multijet background, added in quadrature. Any data-to-prediction disagreement is covered by the $t\bar{t}$ and/or W +jets normalisation and modelling uncertainties detailed in Section 10. In Figure 5(a) and Figure 6(a), the origin of the well-modelled bumps around $|\eta| = 2.5$ is the same as for Figure 3(a). In addition, the well-modelled decrease at $|\eta| = 2$ shown in Figure 5(a) is due to the rejected events in the W +jets control region, which satisfy the signal selection requirement of $|\eta(\text{non } b\text{-jet})| > 2.0$.

9 Analysis of angular distributions

The model introduced in Section 2 is based on the angles θ , θ^* and ϕ^* . The distributions of these angular observables, for events satisfying the signal selection criteria, are shown in Figure 7. Isolation requirements placed on the leptons influence the shape of these angular distributions. Thus from Figure 2 one can see that for $\cos \theta = -1$, the spectator jet overlaps with the b -tagged jet. Similarly, for $\cos \theta^* = -1$, the lepton overlaps with the b -tagged jet. Therefore, in both cases, the acceptance is significantly reduced. For $\cos \theta = +1$, the acceptance is maximal since the spectator jet and the b -tagged jet are back-to-back. For $\cos \theta^* = +1$, although the lepton and the b -tagged jet are back-to-back, the acceptance is not maximal since the lepton is in the same plane as the spectator jet and therefore it may overlap with this jet. For $\phi^* = 0, \pi$ or 2π , the lepton is in the same plane as the spectator jet and therefore it may overlap with this jet. This is disfavoured by the isolation criteria, so acceptance reduces in these three regions. Acceptance is maximal for $\phi^* = \pm\pi/2$, since the lepton is in a plane perpendicular to the spectator.

Just as the angular distribution for the true signal can be expressed in terms of the angular coefficients, $a_{k,l,m}$, of a finite series of orthonormal functions, the reconstructed angular distribution can be expressed as an infinite series of the same functions, similarly to Eq. (1):

$$\varrho_r(\theta, \theta^*, \phi^*; \vec{\alpha}, P) = \sum_{\kappa, \lambda, \mu} \mathcal{A}_{\kappa, \lambda, \mu}(\vec{\alpha}, P) M_{\kappa, \lambda}^{\mu}(\theta, \theta^*, \phi^*), \quad (3)$$

where $|\mu| \leq \min(\kappa, \lambda)$. Multiplying Eq. (3) by $M_{\kappa, \lambda}^{\mu*}(\theta, \theta^*, \phi^*)$, integrating, and applying the orthonormality of the M -functions, one projects out the angular coefficients, obtaining

$$\mathcal{A}_{\kappa, \lambda, \mu} = \int \varrho_r(\theta, \theta^*, \phi^*; \vec{\alpha}, P) M_{\kappa, \lambda}^{\mu*}(\theta, \theta^*, \phi^*) d(\cos \theta) d\Omega^*.$$

For a discrete set of data that follows ϱ_r , the angular coefficients can be estimated as the average value of the function over the data:

$$\mathcal{A}_{\kappa, \lambda, \mu} = \langle M_{\kappa, \lambda}^{\mu*}(\theta, \theta^*, \phi^*) \rangle,$$

similar to a MC estimation of an integral. Experimental values of these coefficients can thus be obtained by taking this average over a set of discrete data for terms up to a maximum κ and λ , determined by the precision of the data. A similar approach to sequential decays is suggested in Ref. [97]. This technique, called orthogonal series density estimation (OSDE) [98], is essentially a Fourier technique to determine moments of the angular distribution. Since $\mathcal{A}_{\kappa, \lambda, \mu} = \mathcal{A}_{\kappa, \lambda, -\mu}^*$, the coefficients with $\mu = 0$ are purely real, while those with $\mu \neq 0$ can be represented by the real and imaginary components of $\mathcal{A}_{\kappa, \lambda, |\mu|}$. These sets

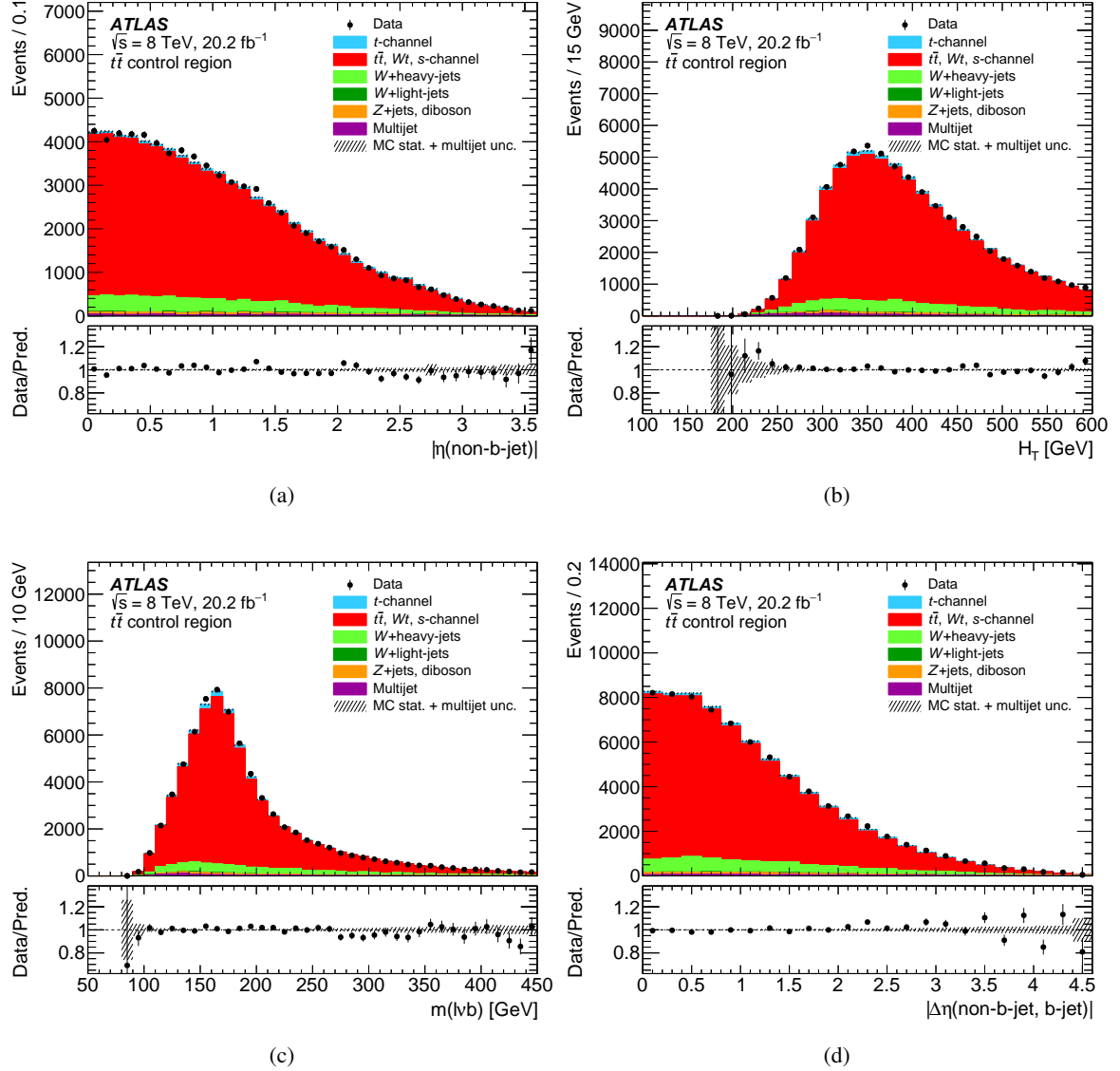


Figure 4: Distributions of (a) $|\eta(\text{non } b\text{-jet})|$, (b) the scalar sum of the p_T of all final-state objects, H_T , (c) reconstructed top-quark mass, $m(\ell\nu b)$, and (d) $|\Delta\eta(\text{non } b\text{-jet}, b\text{-jet})|$ in the $t\bar{t}$ control region for the merged electron and muon channels. The multijet background is estimated using data-driven techniques, while contributions from simulated W +jets, top-quark backgrounds and t -channel event samples are normalised to the results of a maximum-likelihood fit to event yields in the signal and control regions. The uncertainty bands correspond to the uncertainties due to the size of the simulated event samples added in quadrature with the data-driven normalisation uncertainty of 70% estimated for the multijet contribution. The lower plots show the ratio of data to prediction in each bin.

of reconstructed and true angular coefficients, $\mathcal{A}_{k,\lambda,\mu}$ and $a_{k,l,m}$, can be represented by two vectors of coefficients, $\vec{\mathcal{A}}$ and \vec{a} . A covariance matrix, $\mathbf{C} = \text{Cov}(\vec{\mathcal{A}})$, is also determined using OSDE, in the standard way by averaging products of two M -functions.

The background's shape and its covariance matrix are determined through an OSDE analysis of a hybrid sample consisting of background events from simulation samples, and selected data events from

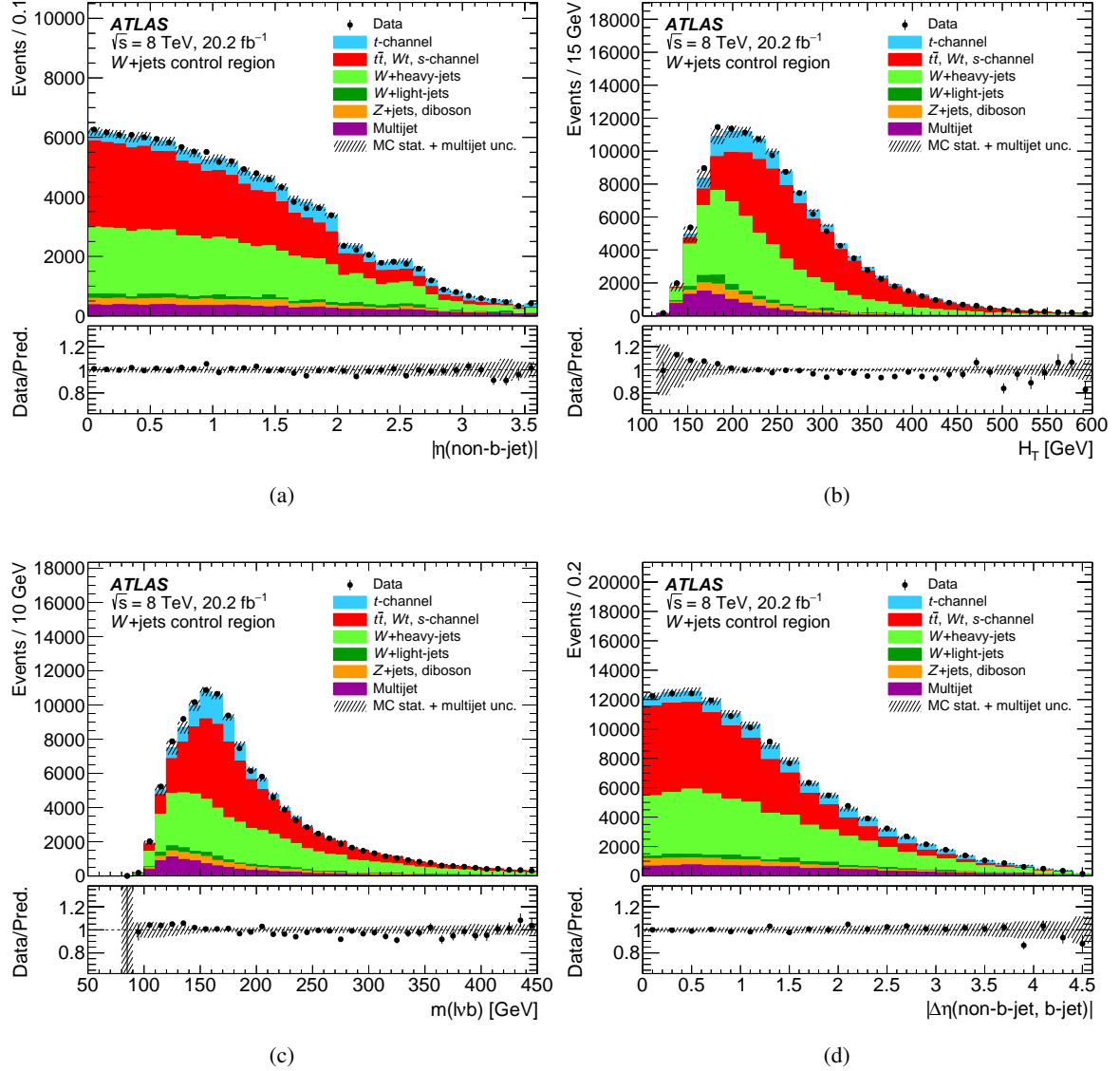


Figure 5: Distributions of (a) $|\eta(\text{non } b\text{-jet})|$, (b) the scalar sum of the p_T of all final-state objects, H_T , (c) reconstructed top-quark mass, $m(\ell\nu b)$, and (d) $|\Delta\eta(\text{non } b\text{-jet}, b\text{-jet})|$ in the W +jets control region for the merged electron and muon channels. The multijet background is estimated using data-driven techniques, while contributions from simulated W +jets, top-quark backgrounds and t -channel event samples are normalised to the results of a maximum-likelihood fit to event yields in the signal and control regions. The uncertainty bands correspond to the uncertainties due to the size of the simulated event samples added in quadrature with the data-driven normalisation uncertainty of 70% estimated for the multijet contribution. The lower plots show the ratio of data to prediction in each bin.

samples enriched in multijet events as reported in Section 7. The vector of reconstructed and background-subtracted coefficients, $\vec{\mathcal{A}}'$, is

$$\vec{\mathcal{A}}' = \frac{1}{f_s} \vec{\mathcal{A}} - \left(\frac{1}{f_s} - 1 \right) \vec{\mathcal{A}}_b,$$

where $\vec{\mathcal{A}}_b$ is the vector of coefficients for the background and f_s is the signal fraction. On the other hand,

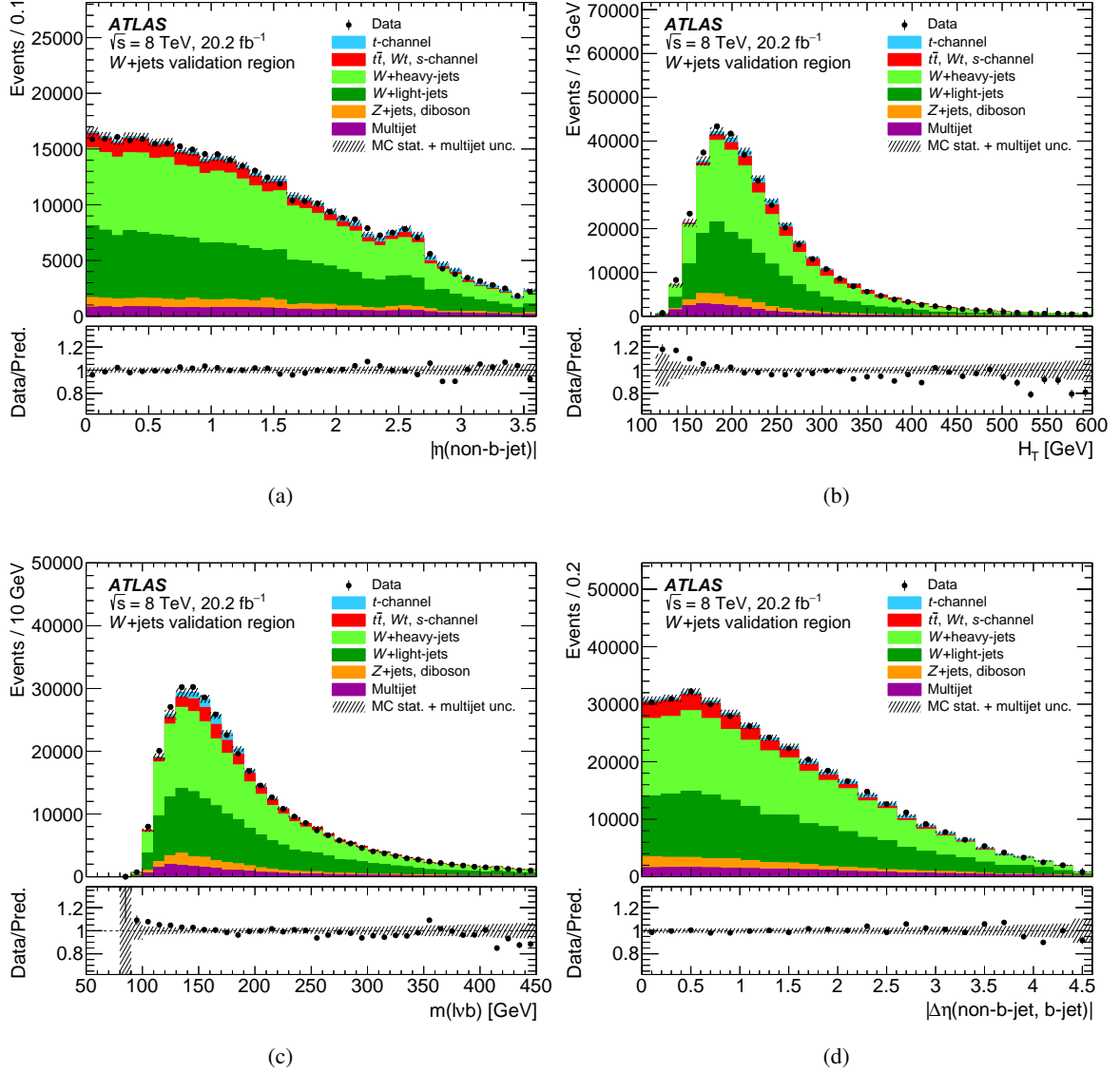


Figure 6: Distributions of (a) $|\eta(\text{non } b\text{-jet})|$, (b) the scalar sum of the p_T of all final-state objects, H_T , (c) reconstructed top-quark mass, $m(\ell\nu b)$, and (d) $|\Delta\eta(\text{non } b\text{-jet}, b\text{-jet})|$ in the W +jets validation region for the merged electron and muon channels. The multijet background is estimated using data-driven techniques, while contributions from simulated W +jets, top-quark backgrounds and t -channel event samples are normalised to the results of a maximum-likelihood fit to event yields in the signal and control regions. The uncertainty bands correspond to the uncertainties due to the size of the simulated event samples added in quadrature with the data-driven normalisation uncertainty of 70% estimated for the multijet contribution. The lower plots show the ratio of data to prediction in each bin.

the covariance matrix \mathbf{C} is modified to include the contribution from the background,

$$\mathbf{C}' = \left(\frac{1}{f_s}\right)^2 \mathbf{C} + \left(\frac{1}{f_s} - 1\right)^2 \mathbf{C}_b, \quad (4)$$

where \mathbf{C}' and \mathbf{C}_b are the covariance matrices of the background-subtracted coefficients and the background coefficients alone, respectively. The second term in Eq. (4) represents a systematic uncertainty in

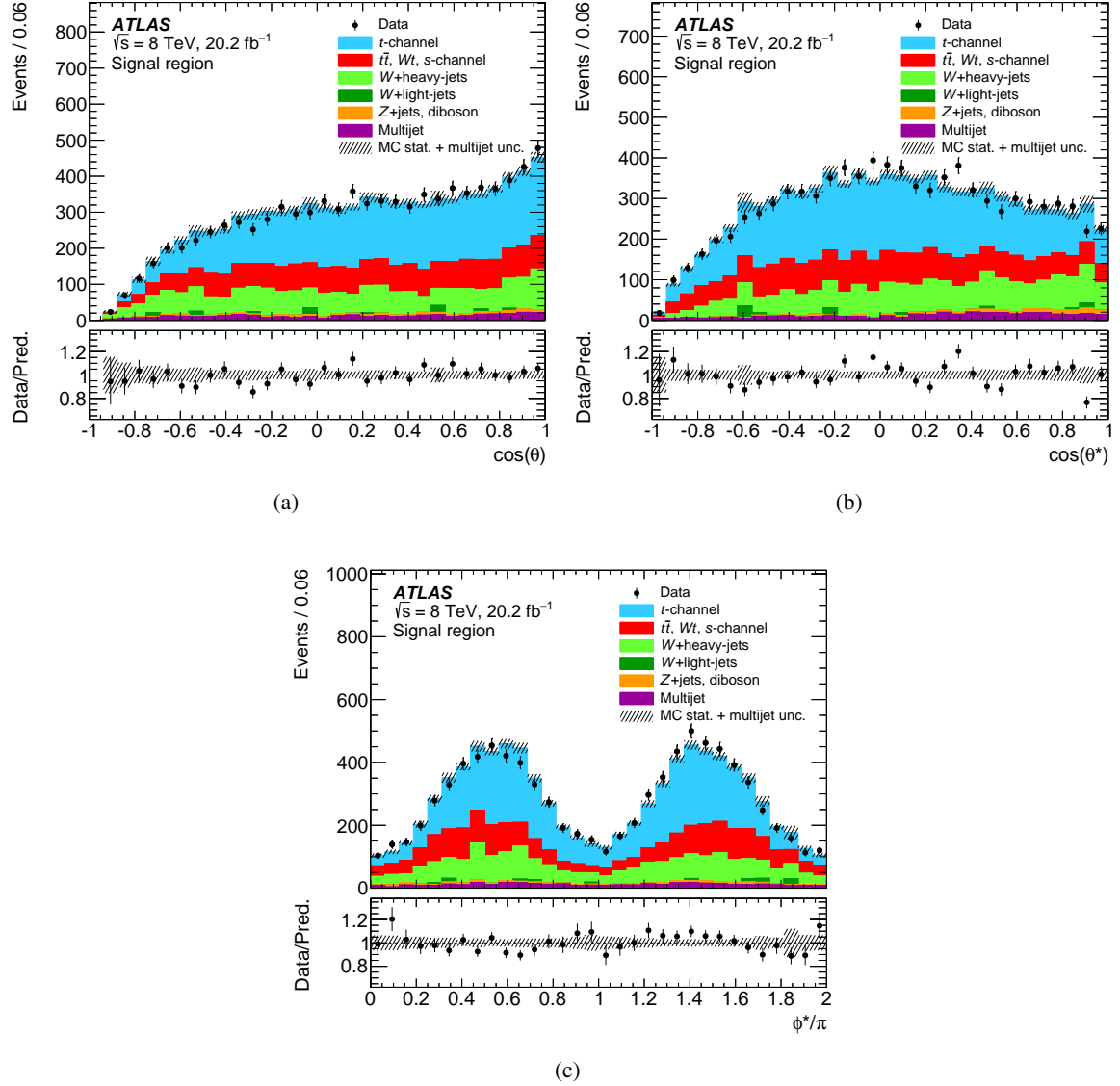


Figure 7: Angular distributions of (a) $\cos \theta$, (b) $\cos \theta^*$ and (c) ϕ^* in the signal region for the electron and muon channels merged, comparing observed data, shown as the black points with statistical uncertainties, to SM signal and background predictions. The multijet background is estimated using data-driven techniques, while contributions from simulated W +jets, top-quark backgrounds and t -channel event samples are normalised to the results of a maximum-likelihood fit to event yields in the signal and control regions. The uncertainty bands correspond to the uncertainties due to the size of the simulated event samples added in quadrature with the data-driven normalisation uncertainty of 70% estimated for the multijet contribution. The lower plots show the ratio of data to prediction in each bin.

\mathbf{C}' due to statistical uncertainties in the background estimate.

Detector effects, both efficiency and resolution, are incorporated through a migration matrix that relates true coefficients, \vec{a} , to reconstructed and background-subtracted coefficients, \vec{A}' . This matrix, denoted by \mathbf{G} , translates all of the nine true coefficients (not counting $a_{0,0,0}$) to the reconstructed coefficients. It is

determined from MC samples produced with the Protos generator using a Fourier analysis of the joint probability density function of true and reconstructed angles, followed by a transformation to coefficients of a conditional probability density function. The procedure is described in more detail in Refs. [36, 37]. In terms of \mathbf{G} ,

$$\vec{\mathcal{A}}' = \mathbf{G} \cdot \vec{a}. \quad (5)$$

Equation (5) cannot be inverted in practice because the matrix \mathbf{G} has more rows than columns, indicating a situation with more equations than unknown variables. Owing to statistical fluctuations or systematic shifts in the measured quantities, it is possible that they cannot all be satisfied simultaneously. The number of rows can be reduced by considering fewer equations. The higher-order terms in $\vec{\mathcal{A}}$ and $\vec{\mathcal{A}}_b$, of which there are an infinite number, are truncated since they represent high-frequency components bringing little information about the true coefficients. In what follows, a truncation is done at $\lambda_{\max} = \kappa_{\max} = 2$ (subscript “max” is the maximum index value of a given series). The maximum values of k and l are chosen to obtain the optimal statistical uncertainty in physics parameters. With this truncation the number of background-subtracted coefficients is 18.

Since a covariance matrix, $\mathbf{C}' = \text{Cov}(\vec{\mathcal{A}}')$, is available, one can minimise the function

$$\chi^2(\vec{a}) = (\vec{\mathcal{A}}' - \mathbf{G} \cdot \vec{a})^T \cdot (\mathbf{C}')^{-1} \cdot (\vec{\mathcal{A}}' - \mathbf{G} \cdot \vec{a}),$$

over the vector \vec{a} . This can be done analytically, and yields the solution

$$\vec{a} = \mathbf{V} \cdot \mathbf{G}^T \cdot (\mathbf{C}')^{-1} \cdot \vec{\mathcal{A}}', \quad (6)$$

with

$$\mathbf{V} = \text{Cov}(\vec{a}) = (\mathbf{G}^T \cdot (\mathbf{C}')^{-1} \cdot \mathbf{G})^{-1}. \quad (7)$$

The deconvolved coefficients, using a migration matrix derived from simulated SM event samples produced with the Protos generator, are shown in Figure 8. Correlations between the different coefficients range from nearly zero to almost 70%. Also shown are the SM predictions, obtained from Eq. (2), using SM values for \vec{a} , and a Protos simulation for the polarisation. Moreover, two new physics scenarios, obtained from Protos simulations, are also shown. The scenario with $\delta_- = \pi$ corresponds to a region where $\text{Re}[g_R/V_L] \approx 0.77$, allowed by the fit in measurements of W boson helicity fractions in top-quark decays [24–28]. The scenario with $f_0^+ = 0.2$ corresponds to a set of couplings ($|V_R/V_L| \approx 0.65$, and $|g_L/V_L| \approx 0.27$) that are also consistent with measurements of W boson helicity fractions, but where 20% of the longitudinal W bosons are due to right-handed couplings.

The derivation of the migration matrix, \mathbf{G} , and background model, $\vec{\mathcal{A}}_b$, described above, is based on the form of these distributions in MC simulation. For the background model, constructed from the sum of all predicted backgrounds with an appreciable effect on the distribution, this includes events containing top quarks, primarily from $t\bar{t}$ production, the distribution of which is affected by changing the values of the anomalous couplings. The efficiency and resolution models are averages over all unmeasured distributions in the signal. Variations in the values of anomalous couplings alter those unmeasured distributions, which could lead to a dependence on these couplings for the efficiency and resolution models. For instance, t -channel single-top-quark production depends on anomalous couplings in both the top-quark production and decay vertices, so varying the couplings alters production-side distributions, such as the p_T and η distributions of the top or spectator quark. Therefore \mathbf{G} and $\vec{\mathcal{A}}_b$ both depend upon \vec{a} . When evaluating \vec{a} for different possible values of \vec{a} , the appropriate values of $\mathbf{G}(\vec{a})$ and $\vec{\mathcal{A}}_b(\vec{a})$ must be used. Consequently, \vec{a} also depends on \vec{a} .

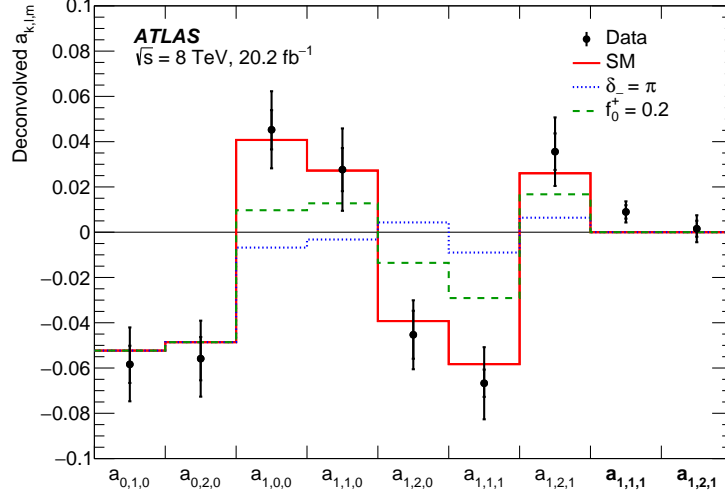


Figure 8: Deconvolved angular coefficients from data using the migration matrix from the SM simulation. Data are shown as black points with statistical uncertainties (inner error bar) and statistical and systematic uncertainties added in quadrature (outer error bar), while SM prediction is shown as a red line. In addition, two new physics scenarios, one with $\delta_- = \pi$ and another one with $f_0^+ = 0.2$, are also shown as a dotted blue line and dashed green line, respectively. The x -axis shows the real and imaginary parts of the angular coefficients, where the latter appears in boldface.

To interpret the measurement of the coefficients $\vec{d}(\vec{\alpha})$ as a measurement of the parameters $\vec{\alpha}$, the real and imaginary parts of the predicted coefficients $a_{k,l,m}$ obtained from Eq. (2) are packed into a vector \vec{d}_{th} . The coefficient $a_{0,0,0}$ is omitted in this procedure because it is constrained by normalisation. Since the number of parameters used to describe the complex coefficients $\dim(\vec{d}) = 9$ exceeds $\dim(\vec{\alpha}) = 6$, an over-constrained system is found. Using $\vec{d}(\vec{\alpha})$ from Eq. (6) and \mathbf{V} from Eq. (7), an additional χ^2 contribution is defined as

$$\chi^2(\vec{\alpha}) = (\vec{d}_{\text{th}}(\vec{\alpha}) - \vec{d}(\vec{\alpha}))^T \cdot \mathbf{V}^{-1} \cdot (\vec{d}_{\text{th}}(\vec{\alpha}) - \vec{d}(\vec{\alpha})) . \quad (8)$$

The final fit uses the combined likelihood

$$-2 \ln \mathcal{L} = \chi^2(\vec{\alpha}) + \chi^2(\vec{d}) . \quad (9)$$

Likelihood profiles over the parameters $\vec{\alpha}$ are computed using a Markov chain MC method [99]. In order to correct for the dependence of \mathbf{G} on $\vec{\alpha}$, the migration matrix is computed on a four-dimensional grid in f_1 , f_1^+ , f_0^+ , and δ_- using Lagrange interpolation between the grid points. Two points are used in f_1^+ , f_0^+ , while four are used in f_1 and δ_- . The range of interpolation is $f_1 \in [0.24, 0.36]$, $f_1^+ \in [0.0, 0.25]$, $f_0^+ \in [0.0, 0.25]$, and $\delta_- \in [-0.5, 0.5]$. The background coefficients $\vec{\mathcal{A}}_b$ are also corrected for the dependence of the $t\bar{t}$ background on $\vec{\alpha}$ in the same manner.

The procedure for deconvolving detector effects has been validated with closure tests, performed using simulation samples produced with the PROTON and ACERMC generators. The model independence of this procedure has been validated using the various simulation samples with anomalous couplings enabled in both the production and the decay vertices, as mentioned in Section 4.

10 Sources of systematic uncertainty

Systematic uncertainties are estimated for the angular coefficients $a_{k,l,m}$. The systematic uncertainties are better behaved in these angular coefficients than in the parameters $\vec{\alpha}$, where they might be close to physical boundaries, e.g. $f_1^+ = 0$ or $f_0^+ = 0$. These systematic uncertainties are used to construct a 9×9 covariance matrix including all correlations between different angular coefficients for each uncertainty considered. The full systematic covariance matrix, \mathbf{V}_{syst} , is then formed by summing the individual matrices. For evaluating the likelihood including the total uncertainty, \mathbf{V}_{syst} is added to the covariance matrix determined from Eq. (7) before evaluating Eq. (8).

Unless addressed specifically, the efficiency and resolution models (i.e. migration matrix) in t -channel events used to estimate the impact of the various sources of uncertainty on the deconvolved measurements are those extracted from the nominal simulation sample produced with the Proros generator and SM couplings. The nominal acceptance and template shape of the t -channel signal is predicted using the Powheg-Box generator. Various signal and background models are determined from MC simulation samples with either alternative generators or parameters varied by their uncertainty in order to estimate systematic uncertainty from different sources. For each source, a likelihood is constructed from the resulting background-subtracted-data model, using events generated with varied parameters. The difference is calculated between the central values estimated at the nominal value of a parameter and at the value varied by its uncertainty, or half the difference between central values estimated with the parameter varied up and down by its uncertainty. These differences are used to construct a covariance matrix for each source of systematic uncertainty. The total covariance matrix for the systematic uncertainties and its correlation matrix are found from the sum of the covariance matrices determined for individual uncertainties.

When estimating the impact of the various sources of uncertainty, the variations are propagated in a correlated way to the rates and to the shapes. The variations due to the systematic uncertainties are also propagated in a correlated way to the signal region and to the two control regions used to constrain the top-quark and W +jets background contributions. For the statistical uncertainties, the variations in the signal and control regions are considered as independent. A set of overall scale factors associated with the top-quark and W +jets backgrounds and with the signal events are extracted for each source of systematic or statistical variation, through the procedure explained in Section 7. The background normalisation is obtained for each systematic uncertainty shift before being subtracted from the observed data. Then the systematic and statistical uncertainties in the fitted normalisation factors are propagated to the measurement.

The sources of systematic uncertainty are split into the following categories:

Detector modelling: The systematic uncertainties in the reconstruction, and energy calibration of electrons and jets and momentum calibration of muons are propagated in the analysis through variations in the modelling of the detector response. Uncertainties related to leptons come from trigger, identification and isolation efficiencies, as well as from the energy or momentum scale and resolution [68, 69]. For jets, the main source of uncertainty is the jet energy scale (JES), evaluated using a combination of in situ techniques [74]. Other jet-related uncertainty sources are the modelling of the energy resolution [100] and reconstruction efficiency [74], the JVF efficiency [75], and the modelling of the tagging efficiencies of b -quark jets, c -quark jets and light-quark jets [77, 78]. The uncertainties from the energy or momentum scale and resolution corrections applied to leptons and jets are propagated to the computation of the $E_{\text{T}}^{\text{miss}}$. The scale and resolution uncertainties due to soft jets and to contributions of calorimeter energy deposits

not associated with any reconstructed objects are also considered independently. For all detector modelling uncertainties, positive and negative uncertainties are estimated separately from the corresponding shifts.

Background normalisation: The uncertainties in the normalisation of the top-quark and W +jets background processes are determined from the scale factor obtained from the maximum-likelihood fit to data. For the top-quark background processes, the statistical post-fit uncertainty of 1% in its overall scale factor is considered. For the W +jets background process, the difference between its nominal overall scale factor and the one estimated when constraining the scale factor of the t -channel contribution to 1.0 in the maximum-likelihood fit (3%) is considered. For the Z +jets and diboson processes, a normalisation uncertainty of 34% is applied to the predictions. For the data-driven normalisation of the multijet background the uncertainty of 70% estimated from the comparison of the matrix method estimates with those given by the jet-electron and anti-muon methods is used. The uncertainty in the integrated luminosity is 1.9% [46] and it is propagated through the normalisation of the simulated background events.

Signal and background modelling: Systematic uncertainties associated with the signal and background modelling are estimated by comparing different generators and by varying parameters in the event generation. The uncertainty in the predicted efficiency and resolution models for the t -channel single-top-quark process, used to deconvolve reconstructed quantities (from PowHEG-Box interfaced to PyTHIA), is estimated by comparing the nominal PROPOS with ACERMC, both interfaced to PyTHIA. This uncertainty also accounts for the difference between models which consider the 4FS in PROPOS and the 5FS+4FS in ACERMC. The uncertainty in the ME calculation in the simulation of the t -channel process is estimated in two ways; by comparing PROPOS with PowHEG-Box, both interfaced to PyTHIA, to account for the mis-modelling of an NLO process by a LO generator, and by comparing PowHEG-Box with MG5_AMC@NLO (ver. 2.2.2) [101], both interfaced to HERWIG (ver. 6.5.20.2) [102] using ATLAS underlying event tune 2 (AUET2) [103], to account for modelling differences between NLO generators. For the $t\bar{t}$ process, PowHEG-Box is compared with MC@NLO (version 4.06) [104], both also interfaced to HERWIG using the AUET2 tune. The uncertainty in the PS and hadronisation is estimated by comparing PowHEG-Box interfaced with PyTHIA and HERWIG for both the t -channel and $t\bar{t}$ processes. The uncertainty in the amount of radiation is evaluated for the t -channel and $t\bar{t}$ processes by comparing the nominal samples with the PowHEG-Box samples generated with varied factorisation and renormalisation scales (and different values of the h_{damp} parameter in the case of the $t\bar{t}$ samples), interfaced to PyTHIA with different hadronisation scales or configurations via alternative Perugia sets of tuned parameters (P2012radHi, P2012radLo, P2012mpiHi and P2012loCR) [60]. In this case, the uncertainty is defined by the shift from the nominal measurement. All these signal and background modelling uncertainties are treated as uncorrelated between t -channel and $t\bar{t}$.

The impact of the flavour composition on the modelling of the W +jets distributions is determined by propagating an uncertainty of 50% in the ratio of the $W+bb$ and $W+cc$ contributions. As reported in Section 8, W +light-jets events give a small contribution in the signal region and no associated modelling uncertainty is taken into account. An additional shape modelling uncertainty is considered for the W +jets contribution by applying an event-by-event shape reweighting procedure. This reweighting is derived in the W +jets validation region from the matching to the data (after subtraction of all processes other than W +jets) in the distribution of the p_T of the W boson.

Systematic uncertainties related to the PDF sets are evaluated for all processes, except for the multijet contribution, in a correlated way. The uncertainty is estimated, following a procedure based on the PDF4LHC prescription [87], by calculating a multidimensional envelope of the uncertainties at 68% CL of the CT10, MSTW2008 NLO and NNPDF2.3 [89] PDF sets. Additionally, an uncertainty due to

possible non-linearities in the polarisation, while not statistically significant, is propagated to the final likelihood contours.

The size of simulation samples: The statistical uncertainty due to the size of simulated background event samples enters through the background coefficients and is estimated during the OSDE analysis of simulated background events. It is evaluated by subtracting, in quadrature, the covariance of the deconvolved coefficients with and without the inclusion of the statistical uncertainties from the background. The statistical uncertainty due to the size of simulated signal event samples enters through the migration matrix and is evaluated by subdividing the simulated signal event samples into 16 equally-sized subsamples. Migration matrices are computed for each subsample, each one being used to deconvolve the full nominal simulation signal sample. From the extracted values for \vec{d} , a covariance matrix is determined, reflecting the size of the MC samples.

The expected statistical uncertainty due to the size of the data sample is evaluated from pseudoexperiments. The covariance matrix is evaluated for each experiment and the matrices are then averaged. The result is taken to be the expected covariance for the signal. The square root of the diagonal elements are the predicted uncertainties in the coefficients.

Table 2 shows the contribution of each source of systematic uncertainty to the most sensitive helicity parameters and coupling ratios. The total systematic uncertainty is obtained by adding in quadrature all the individual systematic uncertainties and the MC statistics uncertainties. Finally, the total statistical and systematic uncertainty is computed by adding all contributions in quadrature.

Source	Helicity parameters		Coupling ratios	
	$\sigma(f_1)$	$\sigma(\delta_-)/\pi$	$\sigma(\text{Re}[g_R/V_L])$	$\sigma(\text{Im}[g_R/V_L])$
Statistical	0.022	0.013	0.030	0.027
Jets	0.029	0.007	0.039	0.009
Leptons	0.014	0.002	0.017	<0.001
E_T^{miss}	<0.001	<0.001	<0.001	<0.001
Generator	0.027	0.006	0.030	0.010
Parton shower and hadronisation	0.004	0.003	<0.001	0.003
PDF variations	0.008	0.004	<0.001	<0.001
Background normalisation	<0.001	<0.001	<0.001	<0.001
Multijet normalisation	<0.001	<0.001	<0.001	<0.001
W +jets shape	0.015	0.005	0.007	0.009
Luminosity	<0.001	<0.001	<0.001	<0.001
MC sample sizes	0.009	0.006	<0.001	0.013
Other	<0.001	<0.001	<0.001	<0.001
Total systematic uncertainty	0.044	0.010	0.061	0.017
Total	0.049	0.017	0.068	0.032

Table 2: Statistical and systematic uncertainties in the measurement of helicity parameters f_1 and δ_- , and of coupling ratios $\text{Re}[g_R/V_L]$ and $\text{Im}[g_R/V_L]$. Uncertainties from individual sources are estimated separately for shifts up and down, and symmetrised uncertainties $\sigma(f_1)$ and $\sigma(\delta_-)$, and $\sigma(\text{Re}[g_R/V_L])$ and $\sigma(\text{Im}[g_R/V_L])$ are given. The statistical uncertainty is calculated by evaluating the likelihood including only the covariance matrix, \mathbf{V} , arising from the data statistics. The total uncertainty is calculated by including \mathbf{V}_{syst} in the likelihood calculation as well as \mathbf{V} . Finally, the total systematic uncertainty is computed by subtracting in quadrature the statistical uncertainty from the total uncertainty.

The leading systematic uncertainties for f_1 come from the jet measurements and the generator modelling. For this parameter, the size of the data sample is also an important source of uncertainty. In the case of δ_- , the leading systematic uncertainties are jet measurements, the generator modelling and MC sample sizes. The measurement of δ_- is dominated by the statistical uncertainty in the data. The leading systematic uncertainties for $\text{Re}[g_R/V_L]$ and $\text{Im}[g_R/V_L]$ are the same as for f_1 and δ_- , respectively.

11 Results

In this section, measurements, limits and distributions obtained from a numerical calculation of the likelihood function (Eq. (9)) are shown in the space of the generalised helicity fractions and phases $\vec{\alpha} \equiv \{f_1, f_1^+, f_0^+, \delta_+, \delta_-\}$ and P , or alternatively of the anomalous couplings $V_{L,R}$, $g_{L,R}$, and P . No external constraints or assumptions are imposed on couplings. Values for parameters of interest can be obtained from likelihood profiles, or joint likelihood contours which show the correlations between the extracted parameters.

Likelihood profiles and a joint likelihood contour for the quantities f_0^+ and f_1^+ are shown in Figure 9. The 68% contours represent the total uncertainty in the measurement.

The limit for f_0^+ , i.e. for the fraction of b -quarks that are right-handed in events with longitudinally polarised W bosons, is

$$\begin{aligned} f_0^+ &< 0.041 & (68\% \text{ CL}), \\ f_0^+ &< 0.085 & (95\% \text{ CL}), \end{aligned}$$

compared with the SM expectation of $f_0^+ = 6 \cdot 10^{-5}$. The limit for f_1^+ , i.e. for the fraction of transversely polarised W boson decays that are right-handed, is

$$\begin{aligned} f_1^+ &< 0.053 & (68\% \text{ CL}), \\ f_1^+ &< 0.120 & (95\% \text{ CL}), \end{aligned}$$

compared with the SM expectation $f_1^+ = 0.001$.

The limits obtained for f_1^+ in this analysis are comparable and complementary to those determined from F_R [24–28], since $F_R = f_1 f_1^+$. However, the quantity f_0^+ is not accessible in measurements of the W boson helicity fractions, as those analyses extract F_0 , which only measures the sum of the contributions of both longitudinal amplitudes. The contributions can only be separated in an analysis with polarised top quarks. Since f_1^+ and f_0^+ are found to be very small, there is no sensitivity to the relative phase δ_+ .

The likelihood profile for the top-quark polarisation P is also obtained and it is shown in Figure 10. This leads to the following constraint on the top-quark polarisation:

$$\begin{aligned} P &> 0.86 & (68\% \text{ CL}), \\ P &> 0.72 & (95\% \text{ CL}). \end{aligned}$$

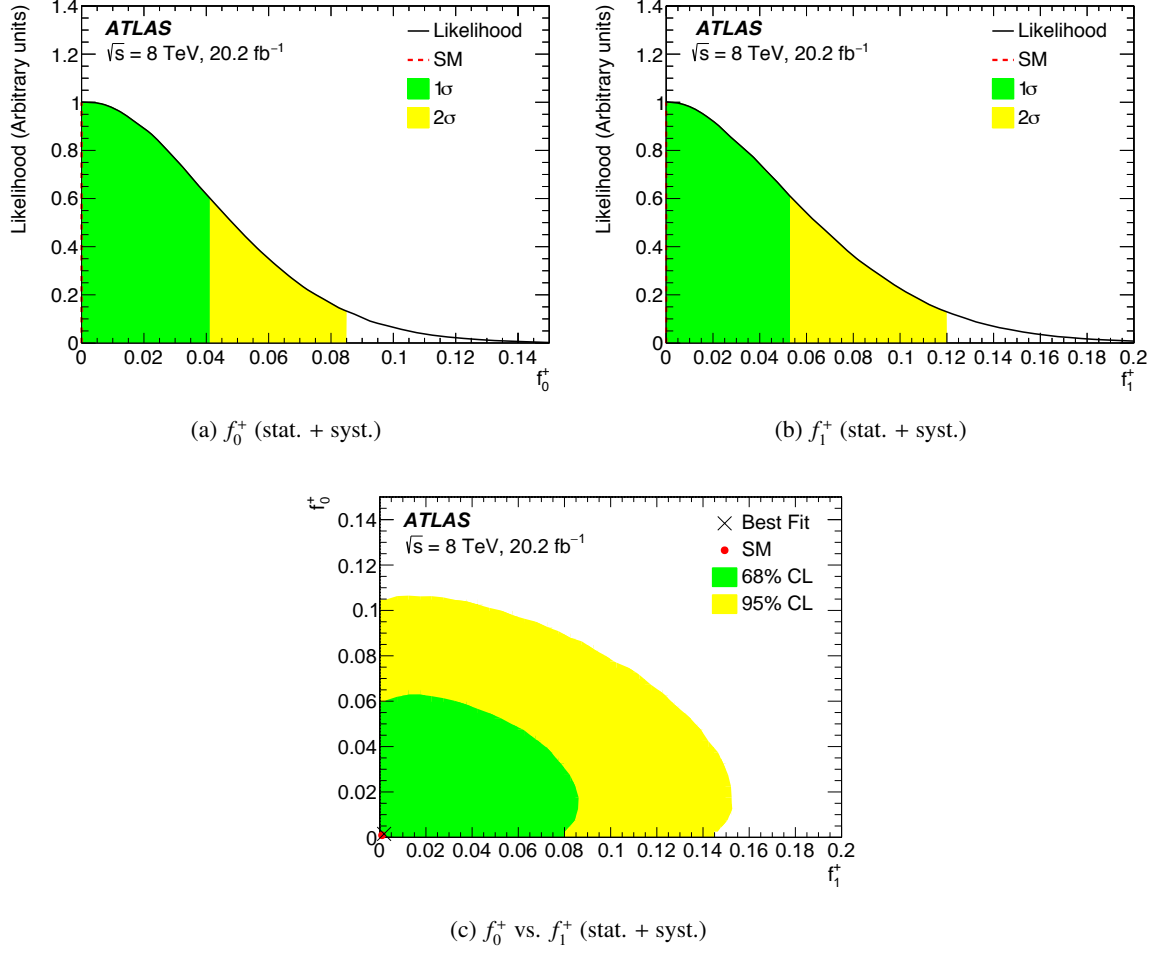


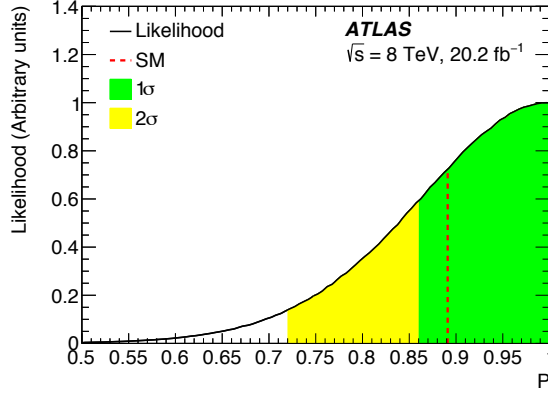
Figure 9: The likelihood profiles for the parameters (a) f_0^+ and (b) f_1^+ are shown. The black line indicates the evaluated likelihood in each bin of the profiled variable. The red dashed line, which overlaps the y -axis, represents the SM expectation. Additionally (c), the joint likelihood profile of f_0^+ as a function of f_1^+ is shown. The red point represents the SM expectation while a black x mark indicates the observed value. Both points overlap with the origin of the x - and y -axis. The 68% and 95% CL regions are shown in green and yellow, respectively.

This is compatible with the SM prediction of $P \approx 0.9$ at $\sqrt{s} = 8$ TeV as computed in Refs. [3, 38, 39], and with recent measurements of the top-quark polarisation obtained from asymmetries of angular distributions with additional inputs on the values of the charged-lepton spin analysing power [105] and/or the W boson helicity fractions [34].

For the parameters for which the analysis obtains point estimates rather than limits, i.e. the fraction f_1 and the phase δ_- as discussed in Section 1, likelihood profiles and a joint likelihood contour are shown in Figure 11. These parameters are measured to be

$$\begin{aligned} f_1 &= 0.296_{-0.023}^{+0.020} \text{ (stat.) } {}_{-0.046}^{+0.043} \text{ (syst.) } = 0.296_{-0.051}^{+0.048}, \\ \delta_- &= 0.002\pi_{-0.014\pi}^{+0.013\pi} \text{ (stat.) } {}_{-0.011\pi}^{+0.010\pi} \text{ (syst.) } = 0.002\pi_{-0.017\pi}^{+0.016\pi}. \end{aligned}$$

Correlations between the coefficients of Figure 8 are taken into account but do not lead to large correlations between these two parameters. The results are compatible with their SM expectations shown



(a) P (stat. + syst.)

Figure 10: The likelihood profile for the top-quark polarisation P is shown. The black line indicates the evaluated likelihood in each bin of the profiled variable. The red dashed line represents the SM expectation. The 68% and 95% CL regions are shown in green and yellow, respectively.

in Section 2, and improve on the measurements from double-differential angular decay rates done at $\sqrt{s} = 7$ TeV by the ATLAS Collaboration [33].

The dependence of the parameters f_1 and δ_- on the top-quark mass is evaluated using t -channel, Wt -channel, s -channel, and $t\bar{t}$ simulation samples with a range of different top-quark masses. A linear dependence is found, resulting from changes in acceptance at different masses, with a slope of -0.005 GeV^{-1} for f_1 and consistent with zero for δ_- . The uncertainty due to the top-quark mass dependence is not included in the total systematic uncertainty since it has a negligible impact on the results.

The results for the generalised helicity fractions and phases can be interpreted in terms of anomalous couplings by propagating the statistical and systematic uncertainties. Although a parameterisation of P in terms of anomalous couplings, obtained from LO MC simulations, exists [106], it is not included in this interpretation. Likelihood profiles and joint likelihood contours for these couplings are shown in Figures 12 and 13. The 68% contours represent the total uncertainty in the measurement. The normalised observables measured in this paper are sensitive to ratios of couplings, which are presented normalised to the dominant coupling in the SM, V_L . The quantities f_1^+ and f_0^+ depend most strongly on two different combinations of V_R and g_L , while the quantities $f_1(1 - f_1^+)$ and δ_- depend more strongly on V_L and g_R . Since the likelihood is determined in terms of all of these quantities simultaneously, no assumptions need to be imposed on couplings in order to produce these distributions. In each case the measured values are consistent with the SM prediction, i.e. $V_R = g_{L,R} = 0$.

The bounds obtained on V_R and g_L are shown in Figure 12. As this analysis yields no constraint on δ_+ , no constraint can be placed on the relative phase between V_R and g_L . Thus, only bounds on the magnitudes,

$$\begin{aligned} |V_R/V_L| &< 0.23 & (68\% \text{ CL}), \\ |V_R/V_L| &< 0.37 & (95\% \text{ CL}), \end{aligned}$$

and

$$|g_L/V_L| < 0.19 \quad (68\% \text{ CL}),$$

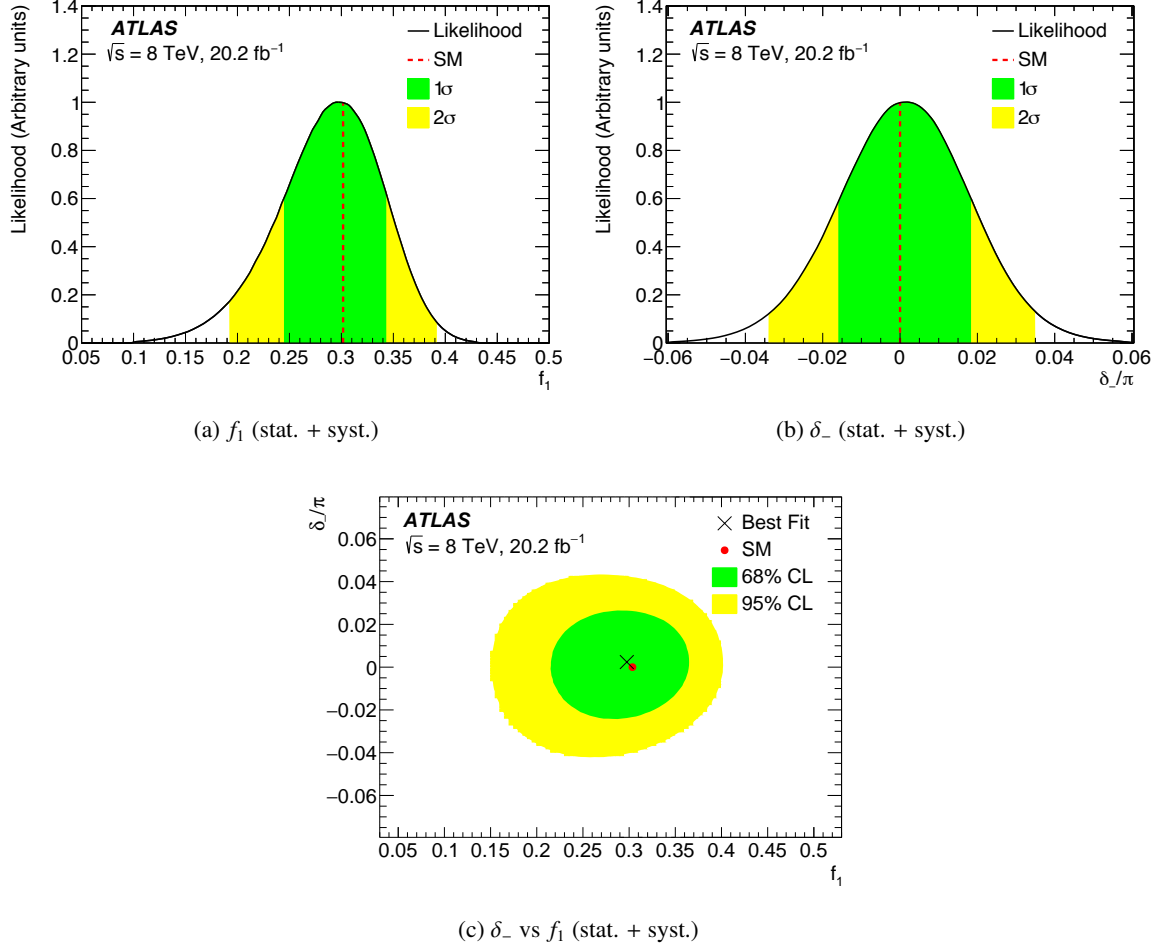


Figure 11: The likelihood profiles for the parameters (a) f_1 and (b) δ_- are shown. The black line indicates the evaluated likelihood in each bin of the profiled variable. The red dashed line represents the SM expectation. Additionally (c), the joint likelihood contour of δ_- as a function of f_1 is shown. The red point represents the SM expectation while a black x mark indicates the observed value. The 68% and 95% CL regions are shown in green and yellow, respectively.

$$|g_L/V_L| < 0.29 \quad (95\% \text{ CL}),$$

are obtained. Limits on these quantities have been obtained from B -meson decays [20], and from measurements of W boson helicity fractions in top-quark decays [24–28], but all of those measurements can only place limits on combinations of couplings, and thus the quoted limits on individual couplings depend on the assumptions made about other couplings.

The propagation of the uncertainties to the $(\text{Re}[g_R/V_L], \text{Im}[g_R/V_L])$ space gives

$$\begin{aligned} \text{Re}\left[\frac{g_R}{V_L}\right] &= 0.006^{+0.033}_{-0.028} (\text{stat.})^{+0.063}_{-0.059} (\text{syst.}) = 0.006^{+0.071}_{-0.065}, \\ \text{Im}\left[\frac{g_R}{V_L}\right] &= -0.005 \pm 0.027 (\text{stat.})^{+0.021}_{-0.012} (\text{syst.}) = -0.005^{+0.034}_{-0.030}. \end{aligned}$$

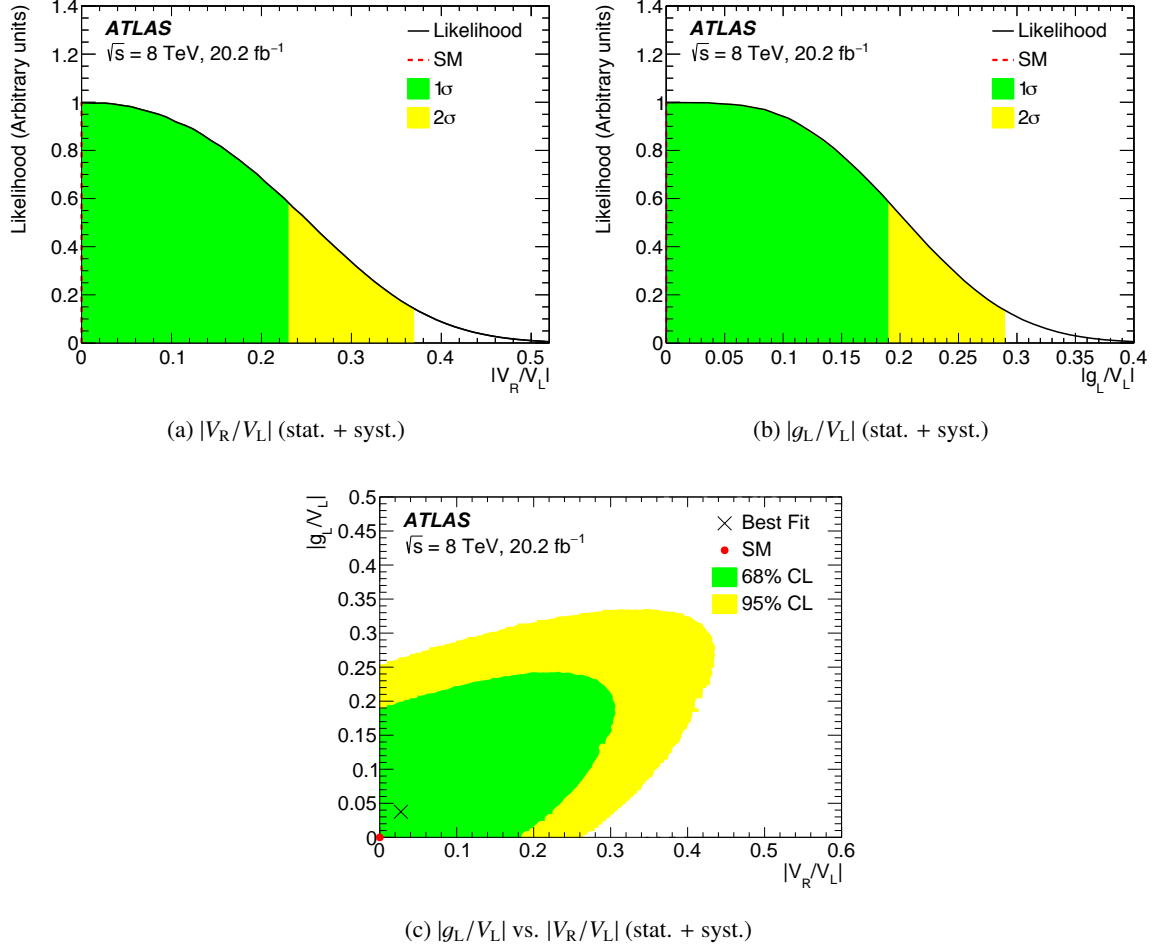


Figure 12: The likelihood profiles for the parameters (a) $|V_R/V_L|$ and (b) $|g_L/V_L|$ are shown. The black line indicates the evaluated likelihood in each bin of the profiled variable. The red dashed line, which overlaps the y -axis, represents the SM expectation. Additionally (c), the joint likelihood contour of $|g_L/V_L|$ as a function of $|V_R/V_L|$ is shown. The red point, which overlaps with the origin of the x - and y -axis, represents the SM expectation while a black x mark indicates the observed value. The 68% and 95% CL regions are shown in green and yellow, respectively.

A linear dependence is found for the coupling ratios on the top-quark mass, which is evaluated with the top-quark mass-varied samples mentioned before. A slope of 0.008 GeV^{-1} is found for $\text{Re}[g_R/V_L]$, while the slope is consistent with zero for $\text{Im}[g_R/V_L]$. Similarly to f_1 and δ_- , the uncertainty due to the top-quark mass dependence is not included in the total systematic uncertainty since it has no significant impact on the results.

Confidence intervals are placed simultaneously on the values of the ratio of the anomalous couplings g_R and V_L at 95% CL,

$$\text{Re}\left[\frac{g_R}{V_L}\right] \in [-0.12, 0.17] \quad \text{and} \quad \text{Im}\left[\frac{g_R}{V_L}\right] \in [-0.07, 0.06].$$

The best constraints on $\text{Re}[g_R]$ derive from measurements of the W boson helicity fractions in top-quark pair decays, with $\text{Re}[g_R] \in [-0.02, 0.06]$ and $[-0.08, 0.07]$, both at 95% CL, from ATLAS at

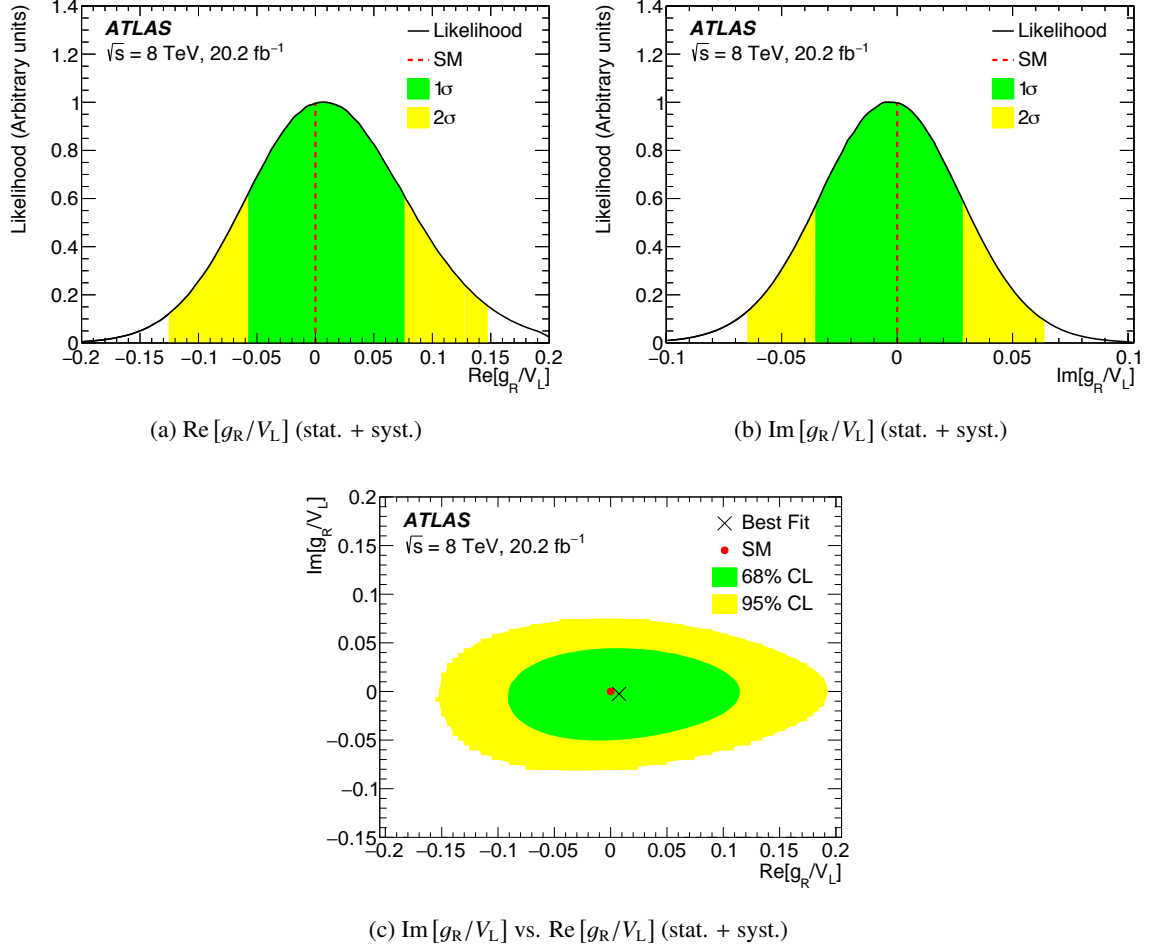


Figure 13: The likelihood profiles for the parameters (a) $\text{Re}[g_R/V_L]$ and (b) $\text{Im}[g_R/V_L]$ are shown. The black line indicates the evaluated likelihood in each bin of the profiled variable. The red dashed line represents the SM expectation. Additionally (c), the joint likelihood contour of $\text{Im}[g_R/V_L]$ as a function of $\text{Re}[g_R/V_L]$ is shown. The red point represents the SM expectation while a black x mark indicates the observed value. The 68% and 95% CL regions are shown in green and yellow, respectively.

$\sqrt{s} = 8$ TeV [26] and from CMS at $\sqrt{s} = 7$ TeV [25], respectively. However, these limits use the measured single-top-quark production cross-section [29, 30] along with the assumption that $V_L = 1$, $\text{Im}[g_R] = 0$, and either $g_L = 0$ or $V_R = 0$. Without these assumptions only a circular region in the complex g_R plane within $0.0 \lesssim \text{Re}[g_R/V_L] \lesssim 0.8$ can be excluded by W boson helicity fractions measurements. The measurements presented here require no assumptions in values of the other anomalous couplings, and on their own can exclude large values of $\text{Re}[g_R/V_L]$.

Along these lines, from the double-differential angular decay rates in t -channel single-top-quark events in ATLAS at $\sqrt{s} = 7$ TeV [33], confidence intervals are placed simultaneously on the coupling ratios, $\text{Re}[g_R/V_L] \in [-0.36, 0.10]$ and $\text{Im}[g_R/V_L] \in [-0.17, 0.23]$, at 95% CL, assuming $V_R = g_L = 0$. Furthermore, slightly better limits on the imaginary part of g_R are set from asymmetries by ATLAS at $\sqrt{s} = 8$ TeV, giving $\text{Im}[g_R] \in [-0.18, 0.06]$ [34], at 95% CL, assuming again $V_R = g_L = 0$. The limits presented in this paper improve on both these results and extend current constraints on g_R to the whole complex plane

by simultaneously measuring information about $\text{Re}[g_R/V_L]$ and $\text{Im}[g_R/V_L]$.

12 Conclusion

The analysis presented in this paper uses the triple-differential decay rate in electroweak production and subsequent decay of single top quarks to constrain the complex parameters of the effective Lagrangian that describes the properties of the Wtb vertex. An analysis of angular distributions of the decay products of single top quarks produced in the t -channel constrains these parameters simultaneously. The analysis is based on 20.2 fb^{-1} of pp collision data at $\sqrt{s} = 8 \text{ TeV}$ collected with the ATLAS detector at the LHC. The selected events contain one isolated electron or muon, large E_T^{miss} , and exactly two jets, with one of them identified as likely to contain a b -hadron. A cut-based analysis is used to discriminate the signal events from background, and the electron and muon channels are merged. An OSDE technique is used to perform an angular analysis of the triple-differential decay rate in order to determine six observables simultaneously, i.e. five generalised helicity fractions and phases, as well as the polarisation of the produced top quark. Detector effects are deconvolved from data using Fourier techniques. The fraction f_1 of decays containing transversely polarised W bosons is measured to be $f_1 = 0.30 \pm 0.05$. The phase δ_- between amplitudes for transversely and longitudinally polarised W bosons recoiling against left-handed b -quarks, is measured to be $\delta_- = 0.002\pi_{-0.017\pi}^{+0.016\pi}$, giving no indication of CP violation. The fractions of transverse and longitudinal W bosons accompanied by right-handed b -quarks are also constrained at 95% CL to $f_1^+ < 0.120$ and $f_0^+ < 0.085$, respectively. The fractions f_1 and f_1^+ are related to the W boson helicity fractions (F_R , F_0 , and F_L), while the fraction f_0^+ , which is previously unmeasured, separates F_0 into two components involving left- and right-handed b -quarks. Based on these measurements, 95% CL intervals are placed on the ratio of the complex coupling parameters g_R and V_L such that $\text{Re}[g_R/V_L] \in [-0.12, 0.17]$ and $\text{Im}[g_R/V_L] \in [-0.07, 0.06]$. Constraints at 95% CL are also placed on the magnitudes of the ratios $|V_R/V_L| < 0.37$ and $|g_L/V_L| < 0.29$, and the polarisation of single top quarks in the t -channel is constrained to be $P > 0.72$ (95% CL). None of the above measurements make assumptions about the value of any of the other parameters or couplings and all of them are in agreement with the SM expectations.

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